

Mathematics for physicists

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Application of Calculus

1. Newton's law of motion in time domain

$$\vec{p}(t) = m \cdot \vec{v}(t) \quad (\text{momentum})$$

$$\vec{v}(t) = \frac{d\vec{r}(t)}{dt} \quad (\text{velocity})$$

$$t \rightarrow \vec{p}(t) ; t + dt \rightarrow \vec{p}(t + dt) = \vec{p}(t) + d\vec{p}(t)$$

$$d\vec{p}(t) = \vec{F}(t) dt \quad (\text{Newton's law of motion})$$

$$\vec{F}(t) \rightarrow \text{Force}$$

in one-dimension case

$$dp_x(t) = F_x(t) dt \quad \text{or} \quad \frac{dp_x(t)}{dt} = F_x(t) \implies m \frac{dv_x(t)}{dt} = m \frac{d^2x(t)}{dt^2} = F_x(t)$$

$$t \rightarrow \vec{r}(t) ; t + dt \rightarrow \vec{r}(t + dt) = \vec{r}(t) + d\vec{r}(t)$$

$$d\vec{r}(t) = ?$$

$$d\vec{r}(t) = \frac{d\vec{r}(t)}{dt} \cdot dt = \vec{v}(t) \cdot dt$$

in one-dim. case

$$dx(t) = \frac{dx(t)}{dt} \cdot dt = v_x(t) \cdot dt$$

$$\frac{d\vec{p}(t)}{dt} = m \frac{d\vec{v}(t)}{dt} = m\vec{a}(t) \rightarrow m\vec{a} = \vec{F}(t) \quad \vec{a} : \text{acceleration}$$

in general

$$\vec{F} = \vec{F}(\vec{r}, \vec{v}) = \vec{F}(\vec{r}(t), \vec{v}(t), \dots)$$

(a) $\vec{F} = -k \cdot \vec{r}$ Hooke's force

one-dim. case $F_x = -k \cdot x$

$$\rightarrow m \frac{d^2 x(t)}{dt^2} = -k \cdot x(t)$$

$$\rightarrow x(t) = A \cdot \cos(2\pi f \cdot t + \delta) \quad (\text{SHM})$$

$$(2\pi f)^2 = \frac{k}{m} \therefore \frac{d^2 \cos(\omega t + \delta)}{dt^2} = -\omega^2 \cos(\omega t + \delta)$$

$$\therefore f = \frac{1}{2\pi} \cdot \left(\frac{k}{m}\right)^{1/2}, \quad x(0) = A \cdot \cos \delta, \quad v(0) = -2\pi f \cdot A \sin \delta$$

$$\therefore \forall x(0), v(0) \implies A, \delta \quad (1)$$

$\therefore A$ and δ are two constants of integration and determined by the initial conditions.

(b) $\vec{F} = -\gamma \cdot \vec{v}$

one-dim. case $F_x = -\gamma \cdot v_x$

$$m \frac{dv_x(t)}{dt} = -\gamma \cdot v_x(t)$$

$$\therefore v_x(t) = B \cdot e^{-\frac{\gamma}{m}t}, \quad B = v_x(0) \quad \therefore \frac{de^{\beta \cdot t}}{dt} = \beta \cdot e^{\beta \cdot t}$$

$$\frac{dx(t)}{dt} = v_x(0) \cdot e^{-\gamma \cdot \frac{t}{m}}$$

$$x(t) = C - \frac{m \cdot v_x(0)}{\gamma} e^{-\gamma \cdot \frac{t}{m}}, \quad x(0) = C - \frac{m \cdot v_x(0)}{\gamma}$$

Note: B and C are two constants of integration and determined by $v_x(0)$ and $x(0)$.

$$\therefore x(t) = x(0) + \frac{m \cdot v_x(0)}{\gamma} \left(1 - e^{-\gamma \cdot \frac{t}{m}}\right)$$

$$x(\infty) = x(0) + \frac{m \cdot v_x(0)}{\gamma}$$

2. Newton's equation of motion in space domain

$$\vec{F} = m\vec{a} = m \frac{d\vec{v}}{dt}$$

one-dim. case $F = m \frac{dv}{dt}$

if $F = F(x)$

then

$$\begin{aligned} F(x) &= m \frac{dv(t)}{dt} = m \frac{dv(x)}{dx} \cdot \frac{dx}{dt} = m \cdot v(x) \cdot \frac{dv(x)}{dx} \\ &= \frac{d\left(\frac{1}{2}mv^2(x)\right)}{dx} \end{aligned}$$

$$\therefore \frac{1}{2}mv^2(x) - \frac{1}{2}mv^2(x_0) = \int_{x_0}^x F(x) \cdot dx = W \quad (\text{work-energy theorem})$$

when $x = x_0 \rightarrow v = v_0 \equiv v(x_0)$

$$\frac{1}{2}mv^2(x) = K \quad (\text{kinetic energy})$$

$$\int_{x_0}^x F(x) \cdot dx = W \quad (\text{work done by } F \text{ from } x_0 \text{ to } x)$$

$$\rightarrow v = v(x) = \frac{dx}{dt} \therefore dt = \frac{dx}{v(x)} \rightarrow t = t(x)$$

(a) Hooke's force $F = -k \cdot x$ if $t = 0$, $x = x_0$, $v = v_0$

$$\frac{1}{2}mv^2 - \frac{1}{2}mv_0^2 = \int_{x_0}^x (-k \cdot x) dx = -\frac{1}{2}k \cdot x^2 + \frac{1}{2}k \cdot x_0^2$$

or

$$\frac{1}{2}mv^2 + \frac{1}{2}k \cdot x^2 = \frac{1}{2}mv_0^2 + \frac{1}{2}k \cdot x_0^2 = E = \text{total mechanical energy}$$

Note $\frac{1}{2}k \cdot x^2 \rightarrow$ elastic potential energy

$$\therefore v = \left(\frac{2}{m} \left(E - \frac{1}{2}k \cdot x^2 \right) \right)^{\frac{1}{2}} = \frac{dx}{dt}$$

$$\therefore dt = \frac{dx}{\left(\frac{2}{m} \left(E - \frac{1}{2}k \cdot x^2 \right) \right)^{\frac{1}{2}}} \implies t - 0 = \int_{x_0}^x \frac{dx}{\left(\frac{2}{m} \left(E - \frac{1}{2}k \cdot x^2 \right) \right)^{\frac{1}{2}}}$$

change of variable: set $x^2 = \frac{2E}{k} \cdot \cos^2 \theta$ i.e. $x = \left(\frac{2E}{k} \right)^{1/2} \cos \theta$, $x = x_0 \implies \theta = \theta_0$

then

$$dx = \left(\frac{2E}{k} \right)^{\frac{1}{2}} (-\sin \theta) d\theta$$

$$\left(\frac{2}{m} \left(E - \frac{1}{2}k \cdot x^2 \right) \right)^{\frac{1}{2}} = \left(\frac{2E}{m} \right)^{\frac{1}{2}} \sin \theta$$

$$\therefore t - 0 = \int_{\theta_0}^{\theta} \left(\frac{m}{k}\right)^{\frac{1}{2}} (-d\theta) = -\left(\frac{m}{k}\right)^{\frac{1}{2}} (\theta - \theta_0)$$

$$\left[x_0 = \left(\frac{2E}{k}\right)^{\frac{1}{2}} \cos \theta_0, t = 0 \implies \theta = \theta_0 \quad (\text{initial phase angle}) \right]$$

$$\therefore \theta - \theta_0 = -\left(\frac{k}{m}\right)^{\frac{1}{2}} \cdot t = -\omega t$$

$$\text{set } \omega \equiv \left(\frac{k}{m}\right)^{\frac{1}{2}} = 2\pi\nu = \text{angular frequency}$$

$$\therefore x = \left(\frac{2E}{k}\right)^{\frac{1}{2}} \cos \theta = \left(\frac{2E}{k}\right)^{\frac{1}{2}} \cos(\omega t - \theta_0)$$

Note: $\omega \cdot T = 2\pi$, $T = \frac{2\pi}{\omega} = \frac{1}{\nu}$; $T \rightarrow$ period, $\nu \rightarrow$ frequency

Note: when $x \rightarrow \max$ i.e. $x = A =$ Amplitude, then $v = 0$

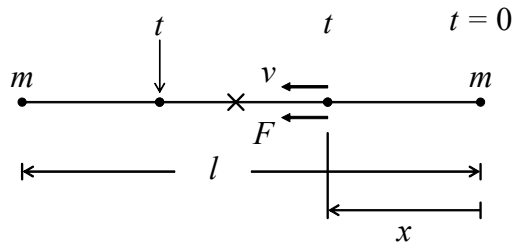
$$\therefore E = \frac{1}{2}k \cdot A^2 \quad A = \left(\frac{2E}{k}\right)^{\frac{1}{2}} \rightarrow x = A \cdot \cos(\omega t - \theta_0)$$

$$\therefore x = A \cdot \cos\left(\sqrt{\frac{k}{m}} \cdot t - \theta_0\right)$$

or

$$t = \left(\frac{m}{k}\right)^{\frac{1}{2}} \left(\cos^{-1} \frac{x_0}{A} - \cos^{-1} \frac{x}{A}\right)$$

(b) Gravitational force



gravitational force acting on the particle at time t :

$$F = \frac{Gm^2}{(\ell - 2x)^2} = F(x)$$

Find the time required for the two initially rest particles to collide.

at time t , $v = \frac{dx}{dt}$, $v(0) = 0$

$$\frac{1}{2}mv^2 - 0 = \int_0^x F(x) \cdot dx = \int_0^x \frac{Gm^2}{(\ell - 2x)^2} \cdot dx$$

$$\begin{aligned} \therefore \frac{1}{2}mv^2 &= \frac{1}{2}Gm^2 \cdot \frac{1}{\ell - 2x} \Big|_0^x = \frac{Gm^2}{2} \left(\frac{1}{\ell - 2x} - \frac{1}{\ell} \right) = Gm^2 \cdot \frac{x}{\ell(\ell - 2x)} \\ \therefore v(x) &= \left(\frac{2Gm}{\ell} \cdot \frac{x}{\ell - 2x} \right)^{\frac{1}{2}} = \frac{dx}{dt} \\ \therefore t_c - 0 &= \int_0^{t_c} dt = \int_0^{\frac{\ell}{2}} \left(\frac{\ell}{2Gm} \cdot \frac{\ell - 2x}{x} \right)^{\frac{1}{2}} \cdot dx \end{aligned}$$

change of variable: set $x = \frac{\ell}{2} \sin^2 \theta$ \therefore $x = 0 \rightarrow \theta = 0 \rightarrow t = 0$
 $x = \frac{\ell}{2} \rightarrow \theta = \frac{\pi}{2} \rightarrow t = t_c$

$$dx = \ell \cdot \sin \theta \cdot \cos \theta \cdot d\theta$$

$$\left(\frac{\ell - 2x}{x} \right)^{\frac{1}{2}} = \left(\frac{\ell \cdot \cos^2 \theta}{\frac{\ell}{2} \sin^2 \theta} \right)^{\frac{1}{2}} = \sqrt{2} \cdot \frac{\cos \theta}{\sin \theta}$$

$$\begin{aligned} \therefore t_c &= \int_0^{\frac{\pi}{2}} \left(\frac{\ell}{Gm} \right)^{\frac{1}{2}} \cdot \ell \cos^2 \theta d\theta = \left(\frac{\ell^3}{Gm} \right)^{\frac{1}{2}} \int_0^{\frac{\pi}{2}} \cos^2 \theta d\theta \\ &= \frac{\pi}{4} \cdot \left(\frac{\ell^3}{Gm} \right)^{\frac{1}{2}} \end{aligned}$$

$$\therefore \int_0^{\pi/2} \cos^2 \theta d\theta = \frac{1}{2} \int_0^{\pi/2} (\cos 2\theta + 1) d\theta = \frac{1}{2} \left(\frac{1}{2} \sin 2\theta + \theta \right) \Big|_0^{\pi/2} = \frac{\pi}{4}$$

Between $t = 0$ and $t = t_c$

$$\begin{aligned} t = t(x) &= \left(\frac{\ell^3}{Gm} \right)^{\frac{1}{2}} \left(\frac{1}{4} \sin 2\theta + \frac{1}{2} \theta \right), \text{ where } \theta = \sin^{-1} \left(\frac{2x}{\ell} \right)^{\frac{1}{2}} \\ \sin 2\theta &= 2 \cdot \left(\frac{2x}{\ell} \right)^{\frac{1}{2}} \cdot \left(1 - \frac{2x}{\ell} \right)^{\frac{1}{2}} \end{aligned}$$

$$\therefore t(x) = \left(\frac{\ell^3}{Gm} \right)^{\frac{1}{2}} \cdot \frac{1}{2} \left[\left(\frac{2x}{\ell} \right)^{\frac{1}{2}} \left(1 - \frac{2x}{\ell} \right)^{\frac{1}{2}} + \sin^{-1} \left(\frac{2x}{\ell} \right)^{\frac{1}{2}} \right]$$