

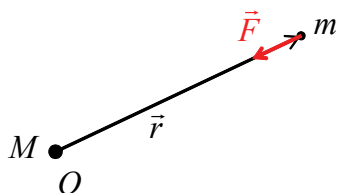
Mathematics for physicists

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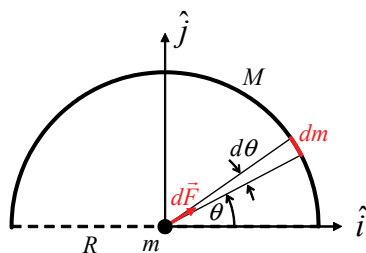
Application of Calculus on Mechanics

1. Universal gravity



$$\vec{F} = -\frac{GMm\hat{r}}{r^2} = -\frac{GMm\vec{r}}{r^3}$$

Example: Total force on m at the center by a uniform semicircular wire with mass M



$$\theta \text{ to } \theta + d\theta \implies dM = \frac{M}{\pi R} \cdot R \cdot d\theta = \frac{M}{\pi} d\theta$$

$$\implies d\vec{F} = dF \cos \theta \cdot \hat{i} + dF \sin \theta \cdot \hat{j}$$

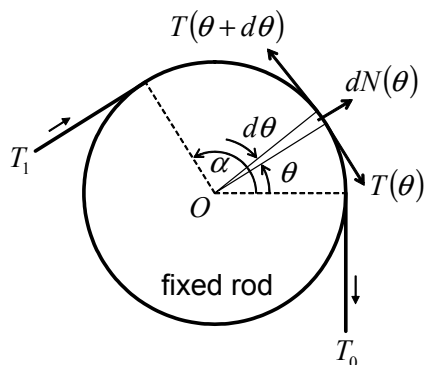
$$dF = \frac{GmdM}{R^2} = G \frac{mM}{\pi R^2} d\theta$$

$$\therefore \vec{F} = \int d\vec{F} = \hat{i} \frac{GMm}{\pi R^2} \underbrace{\int_0^\pi \cos \theta \cdot d\theta}_{=\sin \theta|_0^\pi=0} + \hat{j} \frac{GMm}{\pi R^2} \underbrace{\int_0^\pi \sin \theta \cdot d\theta}_{=-\cos \theta|_0^\pi=2}$$

$$\vec{F} = \hat{j} \frac{2GMm}{\pi R^2}$$

2. Friction force and Tension: $f_\mu = \mu N$, f_μ : sliding friction force, N : normal force, μ : coefficient of sliding friction

Example: Lift a weight by a string tossed over a fixed circular rod



$$\theta \rightarrow T(\theta) \quad \theta + d\theta \rightarrow T(\theta + d\theta) = T(\theta) + dT(\theta)$$

Normal component:

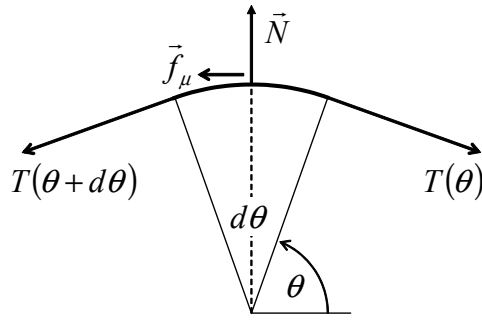
$$(T + dT) \sin \frac{d\theta}{2} + T \sin \frac{d\theta}{2} = dN \quad \text{Normal force}$$

Tangent component:

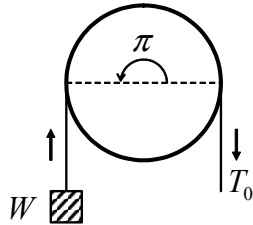
$$(T + dT) \cos \frac{d\theta}{2} - T \cos \frac{d\theta}{2} = -\mu dN \quad \text{Friction force}$$

$$\therefore d\theta \ll 1 \therefore \sin \frac{d\theta}{2} \rightarrow \frac{d\theta}{2}; \quad \cos \frac{d\theta}{2} \rightarrow 1$$

$$\therefore dT \cdot d\theta \ll T \cdot d\theta \quad \text{the } dT \cdot d\theta \text{ term is neglected}$$



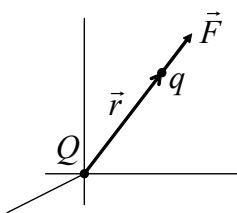
$$\begin{aligned} \therefore T \cdot d\theta &\approx dN \quad \text{and} \quad dT \approx -\mu dN \\ \frac{dT}{T \cdot d\theta} &\approx -\mu_k \quad \text{or} \quad \frac{dT}{d\theta} = -\mu \cdot T \\ \therefore T &= A \cdot e^{-\mu\theta} \quad T(0) = T_0 = A \\ T(\alpha) &= T_1 = T_0 e^{-\mu\alpha} \end{aligned}$$



$$T_0 = W e^{\mu\pi} \quad \therefore T_0 > W$$

Application of Calculus in Electrostatics

3. Coulomb's force

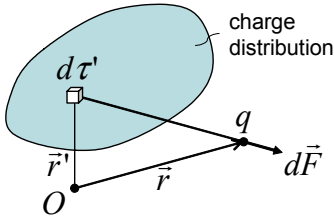


force on point charge q by Q at origin

$$\vec{F} = k \cdot \frac{Q \cdot q \cdot \hat{r}}{r^2}$$

$$\vec{r} = r \cdot \hat{r} \quad \therefore \hat{r} = \vec{r}/r$$

$$\vec{F} = k \frac{Q \cdot q \cdot \vec{r}}{r^3}$$

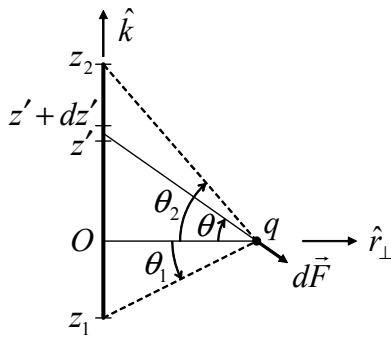


$dQ = \rho(\vec{r}') d\tau'$, $\rho(\vec{r}')$: charge density, $d\tau'$: space element

$$\vec{F} = \int d\vec{F}$$

$$d\vec{F} = k \cdot \frac{\rho(\vec{r}') d\tau' \cdot q \cdot (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

Example: Total Coulomb force on q by a uniformly charged line:



$$\rho_L = \frac{dQ}{dz} \quad \text{linear charge density, } dQ = \rho_L \cdot dz'$$

$$\vec{F} = \int d\vec{F}$$

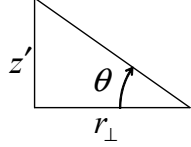
$$d\vec{F} = dF \cos \theta \cdot \hat{r}_\perp + dF \sin \theta \cdot (-\hat{k})$$

$$\therefore \vec{F} = F_{r_\perp} \cdot \hat{r}_\perp + F_k \cdot \hat{k}$$

$$dF = k \cdot \frac{\rho_L dz' \cdot q}{z'^2 + r_\perp^2}$$

$$F_{r_\perp} = \int dF \cdot \cos \theta = \int k \cdot \frac{\rho_L dz' \cdot q}{z'^2 + r_\perp^2} \cdot \cos \theta$$

change of variable from z' to θ :



$$\begin{aligned}
 z' &= r_{\perp} \cdot \tan \theta \\
 \therefore z'^2 + r_{\perp}^2 &= r_{\perp}^2 \cdot \sec^2 \theta \\
 dz' &= r_{\perp} \cdot \sec^2 \theta d\theta \quad \left(\because \frac{d \tan \theta}{d\theta} = \sec^2 \theta \right)
 \end{aligned}$$

$$\therefore F_{r_{\perp}} = k \cdot \rho_L \cdot q \frac{r_{\perp}}{r_{\perp}^2} \int \frac{\sec^2 d\theta}{\sec^2 \theta} \cdot \cos \theta = \frac{k \rho_L q}{r_{\perp}} \int_{\theta_1}^{\theta_2} \cos \theta d\theta = \frac{k \rho_L q}{r_{\perp}} (\sin \theta_2 - \sin \theta_1)$$

Similarly:

$$F_z = -\frac{k \rho_L q}{r_{\perp}} \int_{\theta_1}^{\theta_2} \sin \theta d\theta = \frac{k \rho_L q}{r_{\perp}} (\cos \theta_2 - \cos \theta_1) = k \rho_L q \cdot \left(\frac{1}{(z_2^2 + r_{\perp}^2)^{1/2}} - \frac{1}{(z_1^2 + r_{\perp}^2)^{1/2}} \right)$$

(1) If $\theta_1 \rightarrow -\frac{\pi}{2}$, $\theta_2 \rightarrow +\frac{\pi}{2} \implies \infty$ -long line charge:

$$F_{r_{\perp}} = \frac{2\pi \rho_L q}{r_{\perp}}, \quad F_z = 0$$

(2) If $\theta_1 \rightarrow 0$, $\theta_2 \rightarrow +\frac{\pi}{2} \implies$ semi- ∞ -long line charge:

$$F_{r_{\perp}} = \frac{k \rho_L q}{r_{\perp}}, \quad F_z = -\frac{k \rho_L q}{r_{\perp}}$$