

Mathematics for physicists

Lecturer: Prof. Ven-Chung Lee

(Dated: June 24, 2006)

1. Forced damped oscillation

1. a) Free oscillation

$$m\ddot{x} = -k \cdot x \quad \rightarrow \quad x(t) = A \cdot \cos(\omega_0 t + \theta_0) \quad , \quad \omega_0^2 = \frac{k}{m} : \text{natural frequency}$$

b) Damped oscillation

$$m\ddot{x} = -k \cdot x - \gamma \cdot \dot{x} \quad \rightarrow \quad x_H(t) = A_1 \cdot e^{\beta_1 \cdot t} + A_2 \cdot e^{\beta_2 \cdot t}$$

$$m\beta^2 + \gamma\beta + k = 0$$

$x_H(t)$ is the transient state solution $x_H(t) \rightarrow 0$ as $t \rightarrow \infty$

c) Forced damped oscillation

$$m\ddot{x} = -k \cdot x - \gamma \cdot \dot{x} + F_{ext}$$

$$F_{ext}(t) = F_0 \cdot \cos \omega \cdot t \quad \rightarrow \quad x_s(t) = B_0(\omega) \cos(\omega t + \delta_0) \quad \text{stable state solution}$$

complete solution $x(t) = x_H(t) + x_s(t)$. $x_s(t)$ is the harmonic response.

Approach (1) Comparison of coefficients: set $x_s(t) = B_1 \cos \omega t + B_2 \sin \omega t$

$$\rightarrow \begin{cases} (-m\omega^2 + k) B_1 + \gamma\omega B_2 = F_0 \\ -\gamma\omega B_1 + (-m\omega^2 + k) B_2 = 0 \end{cases} \implies B_1, B_2 \implies B_0, \delta_0$$

Approach (2) Phasor, complex domain

$$\because \cos \omega t = \operatorname{Re}(e^{i\omega t}) \quad , \quad \operatorname{Re} : \text{Real part of}$$

$$\therefore F_{ext}(t) = F_0 \cdot \cos \omega t = \operatorname{Re}(F_0 \cdot e^{i\omega t}) = \operatorname{Re}(\tilde{F}_{ext})$$

$$x_s(t) = B_0 \cdot \cos(\omega t + \delta) = \operatorname{Re}(B_0 e^{i\omega t} \cdot e^{i\delta}) = \operatorname{Re}(\tilde{x}_s(t))$$

$$\therefore m\ddot{x} + \gamma\dot{x} + kx = F_0 \cdot \cos \omega t \rightarrow \operatorname{Re}(m\ddot{\tilde{x}} + \gamma\dot{\tilde{x}} + k\tilde{x} = F_0 e^{i\omega t})$$

$$\text{set } \tilde{x}_s(t) = B_0 e^{i(\omega t + \delta)}$$

$$\rightarrow (-m\omega^2 + i\omega\gamma + k) \cdot B_0 \cdot e^{i(\omega t + \delta)} = F_0 \cdot e^{i\omega t}$$

$$\therefore B_0 e^{i(\omega t + \delta)} = \frac{F_0 \cdot e^{i\omega t}}{(k - m\omega^2) + i\omega\gamma}$$

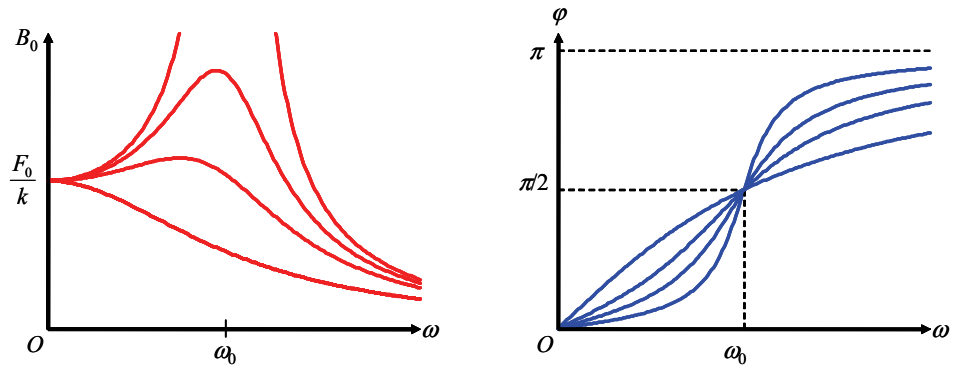
$$\text{set } \tilde{z} = (k - m\omega^2) + i\omega\gamma = \left((k - m\omega^2)^2 + \omega^2\gamma^2 \right)^{\frac{1}{2}} e^{i\varphi} = z e^{i\varphi}$$

$$\tan \varphi = \frac{\omega\gamma}{k - m\omega^2}$$

$$\therefore B_0 = \frac{F_0}{z} \quad \text{and} \quad \delta = -\varphi$$

$$B_0 = \frac{F_0}{\left((k - m\omega^2)^2 + \omega^2\gamma^2 \right)^{\frac{1}{2}}}, \quad B_0 \rightarrow \frac{F_0}{k} \text{ as } \omega \rightarrow 0 \text{ static}$$

$$B_0 \rightarrow 0 \text{ as } \omega \rightarrow \infty, \quad B_0 \rightarrow \frac{F_0}{\gamma \cdot \omega} \text{ as } \omega = \left(\frac{k}{m} \right)^{\frac{1}{2}} = \omega_0 \text{ Resonance}$$

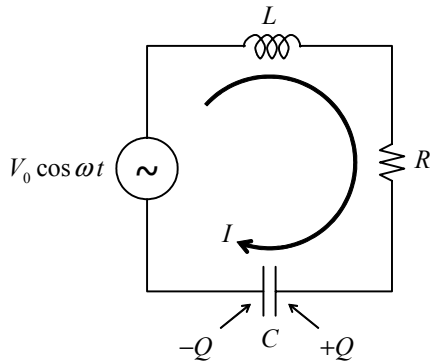


Note: $z \rightarrow \min$ when $dz/d\omega = 0 \rightarrow \omega_r = \left(\frac{k}{m} - \frac{\gamma^2}{2m^2} \right)^{\frac{1}{2}}$ resonance frequency

$$\therefore \omega_r \rightarrow \omega_0 \text{ as } \gamma^2 \ll 2m^2$$

$$x_s(t) = \frac{F_0}{z} \cos(\omega t - \varphi)$$

2. LCR series circuit driven by ACV



$$V_L = L \frac{dI}{dt} \quad \text{voltage of inductance } L$$

$$V_R = R \cdot I \quad \text{Ohm's law}$$

$$V_C = \frac{Q}{C} \quad C : \text{capacitance}$$

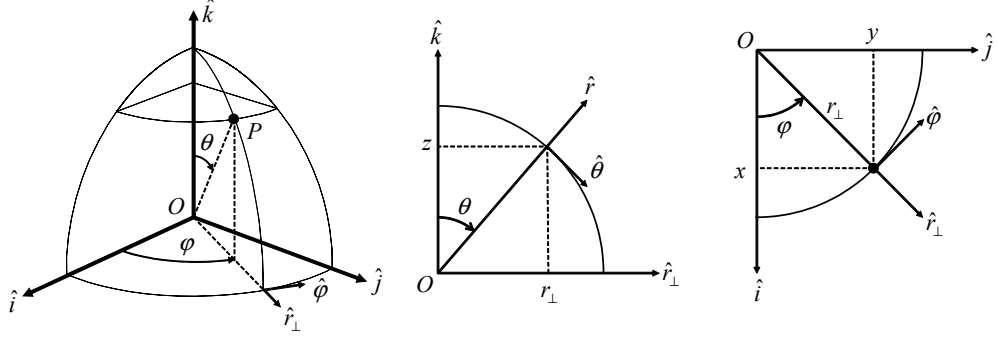
$$V_0 \cos \omega t = V_L + V_C + V_R = L \frac{dI}{dt} + \frac{Q}{C} + R \cdot I$$

$$\text{as } I = \frac{dQ}{dt}$$

$$\therefore V_0 \cos \omega t = L\ddot{Q} + R \cdot \dot{Q} + \frac{1}{C}Q$$

Similar math. as that of the forced damped oscillation.

3. Vector calculus coordinates of $P : (x, y, z), (r, \theta, \varphi), (r_\perp, \varphi, z)$



$$\overrightarrow{OP} = \vec{r} = x\hat{i} + y\hat{j} + z\hat{k} = r \cdot \hat{r} = r_\perp \cdot \hat{r}_\perp + z \cdot \hat{k}$$

$$\text{velocity } \vec{v} = \frac{d\vec{r}}{dt}$$

$$d\vec{r} = d(r \cdot \hat{r}) = dr \cdot \hat{r} + r \cdot d\hat{r}$$

$$d\hat{r} = \hat{r}(r + dr, \theta + d\theta, \varphi + d\varphi) - \hat{r}(r, \theta, \varphi)$$

$$= d\theta \cdot \hat{\theta} + \sin \theta \cdot d\varphi \cdot \hat{\varphi}$$

$$\therefore \vec{v} = \dot{r}\hat{r} + r\dot{\theta}\hat{\theta} + r \cdot \sin \theta \cdot \dot{\varphi}\hat{\varphi} = v_r \cdot \hat{r} + v_\theta \cdot \hat{\theta} + v_\varphi \cdot \hat{\varphi}$$

$$\begin{aligned} \vec{a} = \frac{d\vec{v}}{dt} &= \ddot{r}\hat{r} + \dot{r}\frac{d\hat{r}}{dt} + \dot{r}\dot{\theta}\hat{\theta} + r\ddot{\theta}\hat{\theta} + r\dot{\theta}\frac{d\hat{\theta}}{dt} + \dot{r} \cdot \dot{\varphi} \sin \theta \hat{\varphi} + r \cdot \ddot{\varphi} \sin \theta \cdot \hat{\varphi} \\ &\quad + r \cdot \dot{\varphi} \cdot \dot{\theta} \cos \theta \cdot \hat{\varphi} + r \cdot \dot{\varphi} \sin \theta \frac{d\hat{\varphi}}{dt} \end{aligned}$$

$$\frac{d\hat{r}}{dt} = \dot{\theta}\hat{\theta} + \sin \theta \cdot \dot{\varphi}\hat{\varphi}$$

$$\frac{d\hat{\theta}}{dt} = -\dot{\theta}\hat{r} + \dot{\varphi} \cos \theta \cdot \hat{\varphi}$$

$$\frac{d\hat{\varphi}}{dt} = -\dot{\varphi}\hat{r}_\perp = -\dot{\varphi}(\cos \theta \hat{\theta} + \sin \theta \hat{r})$$

$$\rightarrow \vec{a} = a_r \cdot \hat{r} + a_\theta \cdot \hat{\theta} + a_\varphi \cdot \hat{\varphi} \quad , \quad a_r = \ddot{r} - r\dot{\theta}^2 - r\dot{\varphi}^2 \sin^2 \theta \quad \text{etc.}$$

a) central force problem

$$\vec{F} = F(r) \cdot \hat{r}$$

gravitational force

$$F(r) = -\frac{GMm}{r^2}$$

then

$$ma_r = F(r) \quad \text{and} \quad a_\theta = a_\varphi = 0.$$

Note:

$$a_\theta \neq \frac{dv_\theta}{dt} \quad \because \quad \frac{d\hat{\theta}}{dt} \neq 0, \quad \frac{d\hat{r}}{dt} \neq 0, \quad \text{etc.}$$