

ESR, Zero-Field Splitting, and Magnetic Exchange of Exchange-Coupled Copper(II)–Copper(II) Pairs in Copper(II) Tetraphenylporphyrin *N*-Oxide[#]Fuh-An Yang,[†] Chih-Wei Guo,[†] Yao-Jung Chen,[†] Jyh-Horung Chen,^{*,†} Shin-Shin Wang,[‡] Jo-Yu Tung,[§] Lian-Pin Hwang,^{||} and Shanmugam Elango[⊥]

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The crystal structures of the dimer form of copper(II) tetraphenylporphyrin *N*-oxide, [Cu(tpp-*N*-O)]₂ (**3**-dimer), and zinc(II) tetraphenylporphyrin *N*-oxide, [Zn(tpp-*N*-O)]₂ (**4**-dimer), were established. The geometry at the copper ion in **3**-dimer is essentially square-pyramidal with one oxygen bridge [O(1A)] occupying the apical site, giving a much larger Cu–O bond distance compared to those at the basal plane. The respective Cu···Cu distance and Cu–O–Cu angle in the core of **3**-dimer are 3.987(4) Å and 148.1(3)°. The Zn(1) atom in **4**-dimer has a distorted square-pyramidal [4 + 1] coordination geometry that gives a τ -value of 0.19. The respective Zn···Zn distance and Zn–O–Zn angle in the dimeric unit of **4**-dimer are 4.025(3) Å and 148.1(2)°. The **3**-dimer displays axial X-band electron paramagnetic resonance spectral features ($E_s = 0$) in the powder state at 4 K, giving $g_{||} = 2.51$ ($A_{||s} = (9.6 \pm 0.2) \times 10^{-3} \text{ cm}^{-1}$) and $g_{\perp} = 2.11$ and in the same powder state at 293 K giving $D_s = 0.0731 \text{ cm}^{-1}$ (as derived from $\Delta M_s = 1$ lines) or 0.0743 cm^{-1} (as derived from the $\Delta M_s = 2$ lines). In addition, **3**-dimer displays a $\Delta M_s = 2$ transition at $g = 4.17$ indicating the presence of spin-exchange coupling. The anisotropic exchange interaction ($D_s^{\text{ex}} = 0.132 \text{ cm}^{-1}$) gives the main contribution to D_s in **3**-dimer. The theoretical fit of the susceptibility and effective magnetic moment data of **3**-dimer in the temperature range of 5–300 K gives $2J = 68 \text{ cm}^{-1}$, $g = 2.01$, $p = 0.06$, and a temperature-independent paramagnetism of $10^{-6} \text{ cm}^3 \text{ mol}^{-1}$. This magnetic susceptibility data indicates that the copper(II) ions in **3**-dimer are coupled in a ferromagnetic manner with the ground-spin triplet stabilized by 68 cm^{-1} with regard to the singlet.

Introduction

Metalloporphyrins with a bridged structure between the central metal and one of the four pyrrole nitrogens have drawn much attention in recent times, and so far, bridged metalloporphyrins with a metal–O–N linkage (metal = Ni,^{1,2} Tl³), metal–N–N linkage (metal = Zn,^{4,5,6} Ni,^{7,8} Fe,^{5,9}

Tl,^{7,10,11} Ga^{10,11}), and metal–C–N linkage (metal = Fe¹²) have been reported. Balch and co-workers^{1,2} reported that

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the insertion of nickel(II) into $H_2(OEP-N-O)$ (OEP = dianion of octaethylporphyrin) produced the complex $Ni(OEP-N-O)$ (**1**), in which the metal ions are bound to a nearly planar set of three nitrogen atoms and one oxygen atom. Due to the difficulties in obtaining a crystal suitable for absolute structural assignment by X-ray method, Balch and co-workers² characterized the complex copper(II) octaethylporphyrin *N*-oxide, $Cu(OEP-N-O)$ (**2**), by infrared, UV–vis, and electron spin resonance (ESR) spectroscopy. Judging by both an isotropic ESR spectrum obtained from a toluene solution at 23 °C and an anisotropic spectrum obtained from a 1% sample (**2**) doped in octaethylporphyrin *N*-oxide at 23 °C, they claimed that **2** has a structure analogous to **1** and is a monomeric species. Owing to the failure in measuring the ESR spectrum for **2** in solid phase, they lost the structural information for the same compound in the powder and crystalline phase. Upon replacement of OEP^{2-} with tpp^{2-} (tpp = dianion of *meso*-tetraphenylporphyrin), complex **2** became a new, monomeric compound $Cu(tpp-N-O)$ (**3**). We have recently reported³ the X-ray structure of diamagnetic, mononuclear, six-coordinate and bridged thallium complex of $Tl(tpp-N-O)(OAc)$. It is interesting to figure out the magnetic property for the paramagnetic species **3** in both solid and solution phases. Structural characteristics of binuclear copper complexes, such as dimeric copper(II) species, have been reported in the past.¹³ ESR studies of copper(II) dimers, namely, di(μ -oxo)-bridged copper(II) complexes, had been extensively investigated by Gatteschi et al. and Pibrow and co-workers^{14–16} to elucidate the exchange interactions between the ground and excited states in these structures. Barker et al.¹⁷ reported the magnetic exchange between neighboring copper centers in copper(II) tetraphenylporphyrin $[Cu(tpp)]$, using ESR spectroscopy to detect the low field $\Delta M_s = 2$ transition. Indeed, it is not easy to prepare the dimeric forms of copper(II)–porphyrin, $[Cu(por)]_2$ (por = porphyrin), and consequently, direct investigation of the intermolecular copper–copper interaction, $J(Cu-Cu)$ (or J_{dd}), is difficult in $[Cu(por)]_2$. Moreover, the formation of dimeric species is common in metalloporphyrin π -cation radical systems.^{18,19} A thorough study of magnetic interaction between paramagnetic metal centers and the porphyrin π -cation radical is relatively limited with copper systems such as $Cu(por)^{+}$, the most studied.¹⁸ The ground state of the dimer of $[Cu(por)^{+}]_2$ is determined by three types of isotropic exchange interactions, namely, $J(Cu-Cu)$ (or J_{dd}), $J(\text{radical-radical})$ (or $J_{\pi-\pi}$), and $J(Cu-radical)$

(or J_{dt}).¹⁷ In $Cu(tpp-N-O)$, the molecules could be associated into dimers by secondary bonds between an electronegative oxygen of one monomer and a copper of the other in the solid phase, giving an approximate square pyramidal geometry around copper. The $Cu(tpp-N-O)$ might exist as the dimer form $[Cu(tpp-N-O)]_2$ (**3**-dimer) in solid state. X-ray diffraction data confirms that the presumed monomer of $Cu(tpp-N-O)$ turned out to be a dimer (**3**-dimer) in solid phase. This **3**-dimer is a dimeric form of $[Cu(por)]_2$. The intermolecular $J_{\pi-\pi}$ and intramolecular J_{dt} interactions are neglected; only the intermolecular J_{dd} interaction needs to be considered in magnetic coupling for **3**-dimer. Consequently, the quantity J_{dd} for **3**-dimer has been experimentally measured in the present work. Upon replacement of Cu(II) in **3**-dimer with zinc(II), a new, diamagnetic compound forms, $[Zn(tpp-N-O)]_2$ (**4**-dimer), still with a dimeric structure. The latter complex of **4**-dimer is an analogous diamagnetic metal complex and is used as a diamagnetic correction for **3**-dimer in the solid-state magnetic susceptibility measurements.

To examine possible pathways for magnetic coupling, we have determined the crystal structures of **3**-dimer and **4**-dimer. Furthermore, two complementary techniques are employed to study the electronic nature of **3**-dimer. Studies of temperature dependence of the magnetic susceptibility and of the effective magnetic moment show that the pair of copper(II) atoms are ferromagnetically coupled, giving an $S = 1$ ground state and $S = 0$ level $|2J|$ above this.^{20,21} Application of the Bleaney and Bowers²² equation permits evaluation of $|2J|$ and an average g value for powder samples. Measurement of the ESR spectra arising from **3**-dimer with the $S = 1$ state and application of the spin Hamiltonian (eq 3) permits derivation of the zero-field splitting parameters D_s and E_s . Eaton and co-workers reported that the distance between the copper atoms in the copper porphyrins dimers was determined from the ratio of the intensity of the half-field transitions to the intensity of the allowed transitions.²³ This report is the first unambiguous analysis of zero-field splitting and magnetic exchange interaction in a copper–porphyrin complex. In addition, some of the results of solution and solid-state ESR measurements for **3**-dimer are described.

Experimental Section

[Cu(tpp-N-O)]₂ (3-Dimer). A mixture of $H_2(tpp-N-O)$ (0.063 g, 0.1 mmol) in CH_2Cl_2 (10 mL) and $Cu(OAc)_2$ (0.04 g, 0.22 mmol) in MeOH (5 mL) was refluxed for 1 h.^{24,25} After it was concentrated, the residue was redissolved in CH_2Cl_2 and the solution was dried with anhydrous Na_2SO_4 and filtered. The filtrate was concentrated again and recrystallized from CH_2Cl_2 –MeOH [1:3 (v/v)], affording the purple solids of **3**-dimer (0.05 g, 0.0361 mmol, 36.1%), which

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was again dissolved in CHCl_3 and layered with MeOH to get purple crystals for single-crystal X-ray analysis. FAB-MS, m/z (assignment, rel intensity): 154 ($[\text{NBA} + \text{H}]^+$, 100); 614 ($[\text{H}_2\text{tpp}]^+$, 9.48); 630 ($[\text{H}_2\text{tpp-N-O}]^+$, 3.38); 691 ($[\text{Cu}(\text{tpp-N-O})]^+$, 48.38); 692 ($[\text{Cu}(\text{tpp-N-O}) + \text{H}]^+$, 39.66). UV-vis spectrum, λ (nm) [$\epsilon \times 10^{-3}$, $\text{M}^{-1} \text{cm}^{-1}$] in CH_2Cl_2 : 429 (52.3), 564.0 (4.2), 604 (4.2).

[Zn(tpp-N-O)]₂ (4-Dimer). A mixture of $\text{H}_2(\text{tpp-N-O})$ (0.063 g, 0.1 mmol) and $\text{Zn}(\text{OAc})_2$ (0.04 g, 0.2 mmol) in pyridine (10 mL) was refluxed for 30 min.^{24,25} After the solution was concentrated, the residue was dissolved in CH_2Cl_2 , dried with anhydrous Na_2SO_4 , and the solution was filtered. The filtrate was chromatographed on silica gel (60 g, 70–230 mesh). The major green band eluting with CH_2Cl_2 –pyridine (10:1 (v/v)) was collected and concentrated. Removal of the solvent and recrystallization from pyridine–hexane (1:3 (v/v)) gave purple crystals of 4-dimer (0.4 g, 0.029 mmol, 29%). Compound 4-dimer was redissolved in a trace amount of pyridine to give purple crystals for single-crystal X-ray analysis.

¹H NMR (400 MHz, CDCl_3 , 20 °C): δ 9.02 (s, H_β); 8.98 (d, H_β), ³J(H–H) = 4.4 Hz); 8.82 (d, H_β), ³J(H–H) = 4.4 Hz); 6.90 (s, H_β); 8.45 (s, *o*-H; *o*-H represents *ortho*-phenyl protons); 8.31 (s, *o*-H); 7.80–7.85 (m, *meta* and *para* phenyl protons). FAB-MS, m/z (assignment, rel intensity): 154 ($[\text{NBA} + \text{H}]^+$, 100); 614 ($[\text{H}_2\text{tpp}]^+$, 15.40); 630 ($[\text{H}_2\text{tpp-N-O}]^+$, 9.94); 631 ($[\text{H}_2\text{tpp-N-O} + \text{H}]^+$, 12.50); 692 ($[\text{Zn}(\text{tpp-N-O})]^+$, 53.20); 693 ($[\text{Zn}(\text{tpp-N-O}) + \text{H}]^+$, 44.83). UV-vis spectrum, λ (nm) [$\epsilon \times 10^{-3}$, $\text{M}^{-1} \text{cm}^{-1}$] in CH_2Cl_2 : 420 (63.4), 571 (2.5), 612 (2.5).

Magnetic Susceptibility Measurements. The solid-state magnetic susceptibility were measured under helium on a Quantum Design MPMS5 SQUID susceptometer from 5 to 300 K at a field of 5 kG. The sample was held in a Kel-F bucket. The bucket had been calibrated independently at the same field and temperature. The raw data for 3-dimer were corrected for the molecular diamagnetism. The diamagnetic contribution of the sample was measured from an analogous diamagnetic metal complex, i.e., 4-dimer. The details of the diamagnetic corrections made can be found in ref 21.

Spectroscopy. ESR spectra were measured on an X-band Bruker EMX-10 spectrometer equipped with an Oxford Instruments liquid helium cryostat. Magnetic field values were measured with a digital counter. The X-band resonator was a dual-mode cavity (Bruker ER 4116 DM). Proton NMR spectra were recorded at 599.95 MHz on a Varian Unity Inova-600 spectrometer locked on deuterated solvent and referenced to the solvent peak. Proton NMR is relative to CDCl_3 at $\delta = 7.24$.

The positive-ion fast atom bombardment mass spectrum (FABMS) was obtained in a nitrobenzyl alcohol (NBA) matrix using a JEOL JMS-SX/SX 102A mass spectrometer. UV-vis spectra were recorded at 20 °C on a Hitachi U-3210 spectrophotometer.

Crystallography. Table 1 presents the crystal data as well as other information for 3-dimer·2CHCl₃ and 4-dimer·pyridine. Measurements were taken on a Bruker AXS SMART-1000 diffractometer using monochromatized Mo K α radiation ($\lambda = 0.71073$ Å). Empirical absorption corrections were made for both complexes. The structures were solved by direct methods (SHELXTL-97)²⁶ and refined by the full-matrix least-squares method. All non-hydrogen atoms were refined with anisotropic thermal parameters, whereas all hydrogen atoms except those protons of pyridine for 4-dimer·pyridine were placed in calculated positions and refined

Table 1. Crystal Data for 3-Dimer·2CHCl₃ and 4-Dimer·Pyridine

formula	C ₉₀ H ₅₈ Cl ₆ Cu ₂ N ₈ O ₂	C ₉₃ H ₅₆ N ₉ O ₂ Zn ₂
fw	1623.18	1462.22
space group	$P\bar{1}$	$P\bar{1}$
cryst syst	triclinic	triclinic
<i>a</i> , Å	11.2579(10)	11.2537(8)
<i>b</i> , Å	13.2668(11)	13.1197(9)
<i>c</i> , Å	13.8772(12)	13.8526(10)
α , deg	74.468(2)	73.0700(10)
β , deg	77.592(2)	77.825(2)
γ , deg	67.561(2)	66.5890(10)
<i>V</i> , Å ³	1830.9(3)	1785.0(2)
<i>Z</i>	1	1
<i>F</i> ₀₀₀	830	758
<i>D</i> _{calcd} , g cm ⁻³	1.472	1.365
(Mo K α), mm ⁻¹	0.859	0.732
<i>S</i>	0.954	1.036
cryst size, mm	0.20 × 0.10 × 0.10	0.30 × 0.20 × 0.10
θ , deg	28.37	28.35
<i>T</i> , K	295(2)	294(2)
no. of reflns measd	9080	8830
no. of reflns obsd	4180	5159
<i>R</i> 1 ^a (%)	7.46	7.09
w <i>R</i> 2 ^b (%)	21.10	20.29

$$^a R1 = [\sum |F_o| - |F_c|] / \sum |F_o|. \quad ^b wR2 = \{ \sum [w(F_o^2 - F_c^2)^2] / \sum [w(F_o^2)^2] \}^{1/2}.$$

Table 2. Selected Bond Distances (Å) and Angles (deg) for Compounds 3-Dimer·2CHCl₃ and 4-Dimer·Pyridine

3-Dimer·2CHCl ₃			
Cu(1)–N(1)	1.991(3)	Cu(1)–O(1A)	2.488(7)
Cu(1)–N(3)	1.995(4)	Cu(1)···Cu(1A)	3.987(4)
Cu(1)–N(4)	2.332(4)	Cu(1)–O(1)	1.653(7)
Cu(1)–N(2)	2.130(4)	Cu(1)–O(1')	1.775(4)
		Cu(1)···O(1'A)	4.505(4)
O(1A)–Cu(1)–N(1)	102.0(3)	Cu(1)–O(1)–Cu(1A)	148.1(3)
O(1A)–Cu(1)–N(4)	118.3(3)	O(1)–Cu(1)–N(1)	98.1(3)
O(1A)–Cu(1)–N(3)	94.3(3)	O(1)–Cu(1)–N(4)	150.2(3)
O(1A)–Cu(1)–O(1)	31.9(3)	O(1)–Cu(1)–N(3)	93.9(3)
N(1)–Cu(1)–N(3)	163.39(14)	N(4)–Cu(1)–N(3)	87.12(14)
N(1)–Cu(1)–N(4)	88.07(13)		
4-Dimer·Pyridine			
Zn(1)–N(1)	1.996(3)	Zn(1)···O(1A)	2.512(5)
Zn(1)–N(2)	2.302(3)	Zn(1)···Zn(1A)	4.025(3)
Zn(1)–N(3)	1.994(3)	Zn(1)–O(1)	1.667(5)
Zn(1)–N(4)	2.136(3)	Zn(1)–O(1')	1.774(4)
		Zn(1)···O(1'A)	4.572(4)
Zn(1)–O(1)–Zn(1A)	148.1(2)	O(1A)–Zn(1)–N(1)	92.6(2)
O(1)–Zn(1)–N(1)	93.5(2)	O(1A)–Zn(1)–N(2)	120.3(2)
O(1)–Zn(1)–N(2)	152.2(2)	O(1A)–Zn(1)–N(3)	103.4(2)
O(1)–Zn(1)–N(3)	98.0(2)	N(1)–Zn(1)–N(3)	163.43(12)
O(1)–Zn(1)–O(1A)	31.9(2)	N(2)–Zn(1)–N(3)	87.77(11)
N(1)–Zn(1)–N(2)	87.61(11)		

with a riding model. Compound 4-dimer·pyridine contains pyridine solvent molecules (with partial occupancy) located on inversion and orientationally disordered about it. Table 2 lists selected bond distances and angles for both complexes.

Results and Discussion

Molecular Structures of 3-Dimer and 4-Dimer. The X-ray structures are depicted in Figure 1a for complex 3-dimer and Figure 1b for 4-dimer. The molecular structures of crystalline 3-dimer may be described as a bimolecular unit in which each pair of centrosymmetrically related copper(II) ions, at a rather short distance of Cu(1)···Cu(1A) = 3.987(4) Å, share oxygen atoms, as shown by a dimeric

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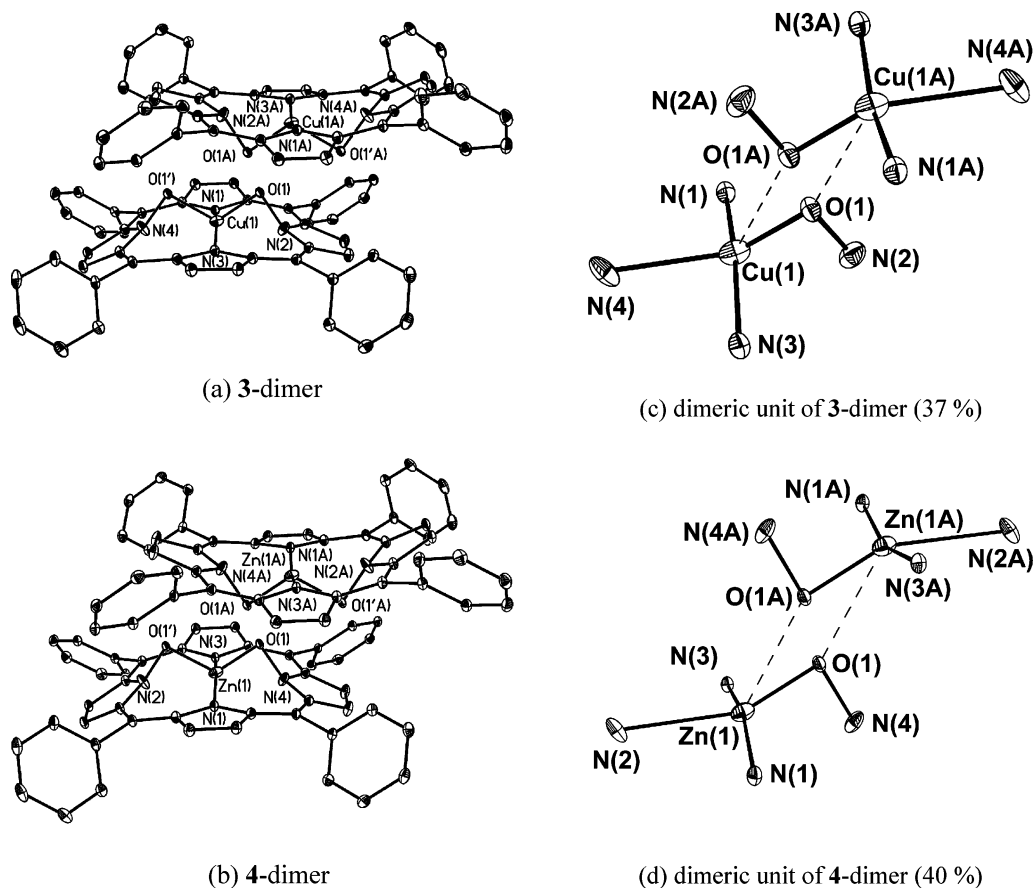


Figure 1. (a) Molecular configuration and atom-labeling scheme for $[\text{Cu}(\text{tpn-}N\text{-O})]_2$ (**3-dimer**) and (b) $[\text{Zn}(\text{tpn-}N\text{-O})]_2$ (**4-dimer**), with ellipsoids drawn at 30% probability. (c) Schematic view of the dimeric copper(II) unit in **3-dimer** (37%). (d) Schematic view of the dimeric zinc(II) unit in **4-dimer** (40%). Hydrogen atoms, solvent CHCl_3 for **3-dimer**, and solvent pyridine for **4-dimer** are omitted for clarity. The oxygen atom is disordered with an occupancy factor of 0.63 (or 0.60) for O(1') and 0.37 (or 0.40) for O(1) within **3-dimer** (or **4-dimer**), respectively.

unit of **3-dimer** (37%) in Figure 1c. Here, the geometry of the coordination of copper(II) is closely related to a distorted square-based pyramid, giving a trigonal distortion parameter (τ) value of 0.22,²⁷ with copper(II) having a coordination of five [4 + 1], CuN_3O_2 , with normal bonds to N(1), N(3), N(4), and O(1) atoms and a longer bond to the O(1A) atom of the centrosymmetrically related molecule (Figure 1c). The Cu(1)···O(1A) bond (2.488(7) Å) is considerably longer than any of the other Cu(1)–N (1.991(3), 1.995(4), and 2.332(4) Å) or Cu(1)–O(1) (1.653(7) Å) bonds. However, the Cu(1)···O(1A) bond length is similar to that of Cu–O = 2.446(5) Å in $\text{Cu}(\text{Pc})(\text{ReO}_4)_2$ ²⁸ but is smaller than those of Cu(1)–O(2) = 2.526(5) Å in $\text{Cu}(\text{Pc})(\text{ReO}_4)_2$ and of Cu(1)–O(1) = 2.533(7) Å in $\text{Cu}_3\text{L}(\text{CH}_3\text{COO})_4$.²⁹ This Cu(1)···O(1A) contact is also significantly shorter than the sum of the van der Waals radii of Cu and O (2.90 Å) and may be described as a weak covalent bond.³⁰ The bond distances (Å) are as follows: Zn(1)···O(1A) = 2.512(5) and Zn(1)–O(1) = 1.667(5), and the mean of Zn(1)–N_p = 2.097(3) for **4-dimer**

(Figure 1b). The Zn(1)···O(1A) distance is longer than the sum of the covalent radii of Zn and O (1.93 Å) but is significantly shorter than sum of the van der Waals radii of Zn and O (2.90 Å).³⁰ This longer Zn(1)···O(1A) contact is too long to be considered as a true coordinated bond and may be viewed as a secondary intermolecular interaction. This kind of secondary interaction was previously observed for $\text{Zn}(\text{tpn})(\text{THF})_2$ with Zn–O = 2.536(7) Å³¹ and for $\text{Zn}[\text{T}(2',6'\text{-DHP})\text{P}](\text{EtOAc})_2 \cdot 2\text{EtOAc}$ with Zn–O = 2.532(3) Å.³² The effective coordination for **4-dimer** is then denoted as [4 + 1], ZnN_3O_2 . The Zn(1) atom in **4-dimer** has a distorted square-pyramidal [4 + 1] coordination geometry that gives a τ -value of 0.19. The Zn(1)···Zn(1A) distance is 4.025(3) Å in the dimeric unit of **4-dimer** (40%) (Figure 1d). The structure suffers from disorder in the location of the oxygen atom. There is only one oxygen atom per **3-monomer** or **4-monomer**. However, the macrocycle of **3-monomer** or **4-monomer** is packed so that the oxygen atom occupies two sites which are shown in Figure 2a with 63% occupancy at O(1') and 37% occupancy at O(1).² The whole molecule of **3-monomer** resulting from these arrangements is shown in parts b and c of Figure 2. A similar kind of disorder with 60% occupancy

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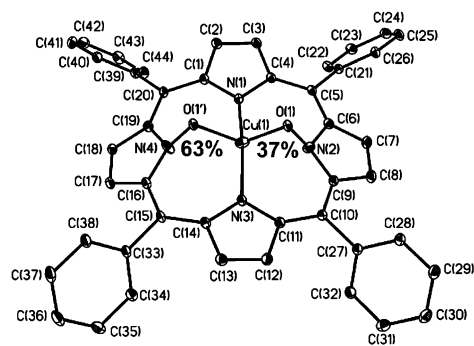
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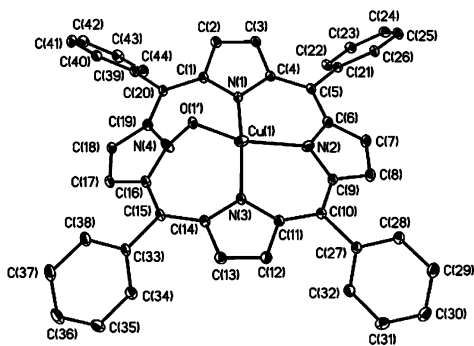
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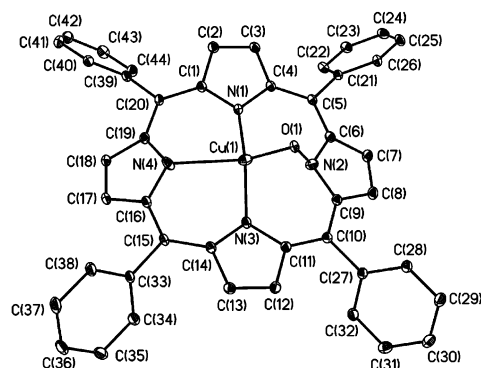
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(a) 3-monomer



(b) 3-monomer (63%)



(c) 3-monomer (37%)

Figure 2. (a) Drawing of the asymmetric unit (3-monomer) in 3-dimer, showing the atom numbering scheme. Due to the disorder, the O(1') and O(1) sites each have 0.63 and 0.37 occupancies, respectively. (b) View of 3-monomer with the site of O(1'), which has an occupancy of 0.63. (c) View of 3-monomer with the site of O(1), which has an occupancy of 0.37.

at O(1') and 40% occupancy at O(1) for 4-monomer is shown at Figure S1 (in the Supporting Information). The M–O values for the major species are shown in Table 2 and in Figure S3 (in the Supporting Information). The geometry around Cu(1) is a distorted square planar in 3-dimer (63%) and that around Zn(1) is also a distorted square planar in 4-dimer (60%) (parts b and f of Figure S3 in the Supporting Information). Bond distances (Å) are Cu(1)–N(1) = 1.991(3), Cu(1)–N(2) = 2.130(4), Cu(1)–N(3) = 1.995(4), and Cu(1)–O(1') = 1.775(4) for 3-dimer (63%); for 4-dimer (60%), the values are Zn(1)–N(3) = 1.994(3), Zn(1)–N(4) = 2.136(3), Zn(1)–N(1) = 1.996(3), and Zn(1)–O(1') =

1.774(4). Since the dimeric unit in the minor occupancy oxygen atom gives a clearly weak axial interaction in the solid state (parts c and d of Figure 1 and parts d and h of Figure S3 in the Supporting Information), all the M–O distances lead to the usage of the minor occupancy oxygen atom.

Spin-Pair Hamiltonian. The Hamiltonian for the exchange coupling between the two copper(II) paramagnetic centers 1 and 2 in 3-dimer with local electronic spin $S_1 = S_2 = 1/2$ can be written as^{33–39}

$$\hat{H} = -2JS_1 \cdot S_2 + S_1 \cdot D_{12} \cdot S_2 + \sum_{i=1}^2 (S_i \cdot D_i \cdot S_i + \beta H g_i \cdot S_i + S_i \cdot A_i \cdot I_i)^{34} \quad (1)$$

where the terms correspond to isotropic exchange coupling, dipolar exchange coupling, single-ion zero-field splitting, electronic Zeeman splitting, and hyperfine interaction ($I_{Cu} = 3/2$) on each copper(II) center, respectively. Here H is the static magnetic field; β is the Bohr magneton; J is the isotropic exchange interaction constant; $A_1, A_2, I_1,$ and I_2 are hyperfine interaction tensors and nuclear spin of ions one and two; $g_1, g_2, D_1,$ and D_2 are the g tensors and the single-ion zero-field splitting tensor of ions one and two, respectively; and D_{12} is the zero-field splitting arising from the coupling of the spins. The latter term is due to anisotropic exchange interactions (D^{ex}) and magnetic dipole–dipole interactions (D^{dip}):

$$D_{12} = D^{ex} + D^{dip} \quad (2)$$

If the term in J is much greater than other terms, the total electronic spin S is a good quantum number and is given by $S = 1, 0$. The Hamiltonian written as eq 1 thus reduces to

$$H_s = -J[S(S+1) - S_1(S_1+1) - S_2(S_2+1)] + \beta H g_s \cdot \vec{S} + D_s \left[S_z^2 - \frac{1}{3}S(S+1) \right] + E_s(S_x^2 - S_y^2) + S \cdot (A_{1,s} \cdot I_1 + A_{2,s} \cdot I_2) \quad (3)$$

Here, D_s and E_s are the axial and rhombic zero-field splitting for the exchange-coupled dimer in the spin state S , respectively; \vec{S} is the total spin operator; and $A_{1,s}$ and $A_{2,s}$ are the hyperfine coupling constants for the S -spin state of nuclei 1 and 2, respectively, in the coupled system.

For spectra of pairs in the strong exchange limit, one can establish the following relations:

$$g_s = \frac{g_1 + g_2}{2} \quad (4)$$

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$$A_{1,s} = \frac{1}{2}A_1; \quad A_{2,s} = \frac{1}{2}A_2 \quad (5)$$

The tensor D_s is related to the component tensors D_1 and D_2 by the relation $D_s = d_1D_1 + d_2D_2 + d_{12}D_{12}$.³⁸ For the $S_1 = 1/2$, $S_2 = 1/2$ system, these coefficients for the spin-coupled system are as follows: $S = 1$, $d_1 = d_2 = 0$, and $d_{12} = 1/2$.

$$D_s = \frac{1}{2}D_{12} = \frac{1}{2}D^{\text{dip}} + \frac{1}{2}D^{\text{ex}} = D_s^{\text{dip}} + D_s^{\text{ex}} \quad (6)$$

D_s is composed of two terms, $D_s^{\text{dip}} = 1/2D^{\text{dip}}$ (from the electron–electron dipolar interaction) and $D_s^{\text{ex}} = 1/2D^{\text{ex}}$ (from spin–orbit coupling) for the coupled system.⁴⁰ D_s^{dip} is related to the Cu–Cu separation r by eq 7^{41,42}

$$D_s^{\text{dip}} = \frac{-(2g_{\parallel}^2 + g_{\perp}^2)\beta^2}{2r^3} \quad (7)$$

ESR Studies. X-band ESR spectra have been obtained for **3**-dimer in powder form at 293 and 4 K (Figure 3) and **3**-monomer in CHCl_3 (or CH_2Cl_2) at 293 K (Figure 4). Figure 3 shows a pair of signals on both sides of the monomer copper species signal with a separation of 780 G and a singlet at about 1 617 G. These two signals are characteristic of triplet dimeric copper species; that is, the paired signals at high field are ascribed to the $\Delta M_s = 1$ transitions and the signal at low field to the $\Delta M_s = 2$ transition. The observation of only a pair of signals at high field suggests that the triplet dimeric species has near or complete axial symmetry with $E_s \approx 0$, the value of 780 G corresponding to a zero-field splitting parameter, D_s .^{43–46} The D_s value derived from the position of the $\Delta M_s = 2$ transition is 0.0743 cm^{-1} ,^{47–49} whereas the value obtained from the $\Delta M_s = 1$ transition is 0.0731 cm^{-1} .⁴⁴ The very good agreement between the two values confirms that the preceding assignment of observed spectrum to a triplet species and the analysis of the fine structure are correct. Lowering the sample temperature to 4 K at perpendicular mode allowed the detection, at high gain, of seven clearly resolved copper hyperfine lines of an average separation of 79.6 G with relative intensities of 1:2:3:4:3:2:1, on the low-field side of the $\Delta M_s = 2$ signal (Figure 3b). This seven-line structures can be attributed to the hyperfine structure from two copper nuclei with $I = 3/2$ and is characteristic of a dimeric Cu(II) triplet state. The similar

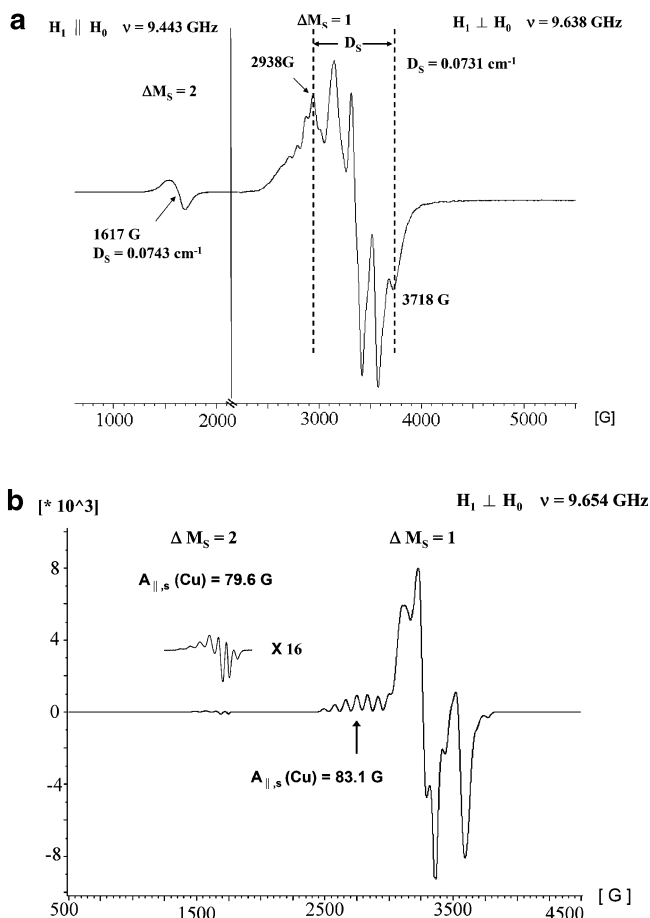


Figure 3. (a) X-band ESR powder spectra for **3**-dimer at 293 K. Low-field region of the spectra: $\Delta M_s = 2$, $H_1 \parallel H_0$, frequency $\nu = 9.443 \text{ GHz}$. High-field region of the spectra: $\Delta M_s = 1$, $H_1 \perp H_0$, frequency $\nu = 9.638 \text{ GHz}$. ESR conditions: microwave power of 20.09 mW ($\Delta M_s = 2$) and 19.97 mW ($\Delta M_s = 1$), magnetic field modulation amplitude of 2 G ($\Delta M_s = 2$) and 1.6 G ($\Delta M_s = 1$), modulation frequency of 100.00 kHz. (b) X-band ESR powder spectra for **3**-dimer at 4 K. ESR conditions for perpendicular mode of spectrum: microwave frequency of 9.654 GHz, microwave power of 2.002 mW, magnetic field modulation amplitude of 1.6 G, and modulation frequency of 100.00 kHz.

kinds of seven-line hyperfine structure were observed with a separation of 74 G (or 81.8 G) in the $\Delta M_s = 2$ transition for dicopper(II) complexes of a 28-membered N_8 macrocycle $[\text{Cu}(\text{HL})(\text{MeCN})_2](\text{BPh}_4)_3$ ⁵⁰ (or $\text{Na}_2\text{Cu-dl-C}_4\text{O}_6\text{H}_2 \cdot 5\text{H}_2\text{O}$).⁵¹ The first six lines with relative intensities of 1:2:3:4:3:2 of a seven-line g_{\parallel} pattern ($g_{\parallel} = 2.51$) are also discernible on the low-field “wing” of the $\Delta M_s = 1$ signal (Figure 3b), and the remainder of the signal is overlapped by the more intense g_{\perp} component ($g_{\perp} = 2.11$). The splitting constant ($A_{\parallel,s}(\text{Cu}) = 83.1 \text{ G}$) in the $\Delta M_s = 1$ transition is close to that with $A_{\parallel}(\text{Cu}) = 77\text{--}82 \text{ G}$ observed for $[\text{Cu}(\text{HL})(\text{MeCN})_2](\text{BPh}_4)_3$ ⁵⁰ and with 78.3 G observed for $\text{Na}_2\text{Cu-dl-C}_4\text{O}_6\text{H}_2 \cdot 5\text{H}_2\text{O}$.⁵¹ The hyperfine coupling constant ($A_{\parallel,s}(\text{Cu}) = 83.1 \text{ G}$) obtained from $\Delta M_s = 1$ is just slightly more than that with ($A_{\parallel,s}(\text{Cu}) = 79.6 \text{ G}$) observed for the $\Delta M_s = 2$ signal, and this average data of $81 \pm 2 \text{ G}$ provides a good estimation of the interaction between two copper centers in **3**-dimer.

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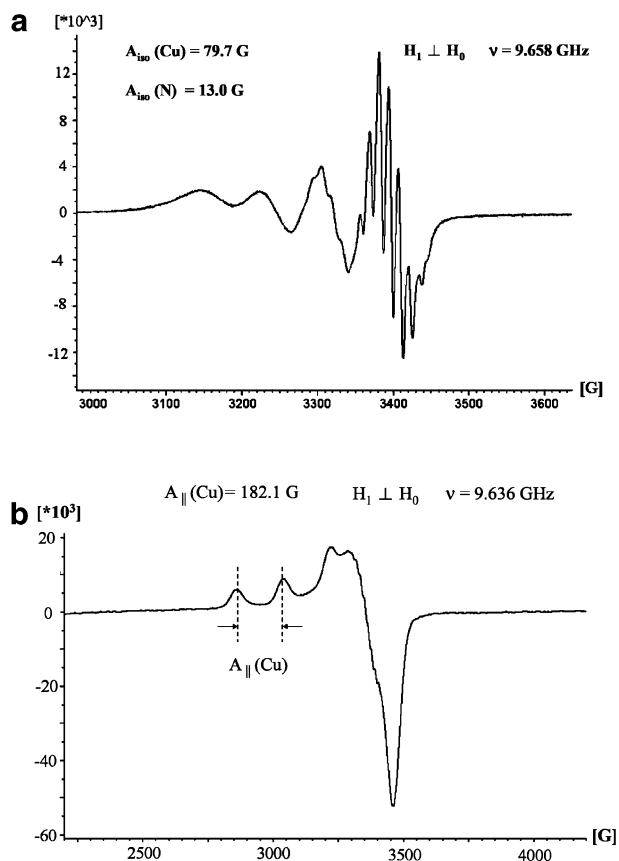


Figure 4. (a) X-band ESR spectra for Cu(tpp-*N*-O) (**3**-monomer) in CHCl₃ (or CH₂Cl₂) at 293 K. ESR conditions: microwave frequency of 9.658 GHz, microwave power of 20.11 mW, modulation amplitude of 2.0 G, and modulation frequency of 100.00 kHz. (b) X-band ESR spectra for Cu(tpp-*N*-O) (**3**-monomer) in CHCl₃ (or CH₂Cl₂) at 77 K. ESR conditions: microwave frequency of 9.636 GHz, microwave power of 20.07 mW, modulation amplitude of 1.6 G, and modulation frequency of 100.00 kHz.

The ESR spectra of **3**-dimer in CH₂Cl₂ at 20 °C shown in Figure 4 indicate that dissociation occurs (since only monomeric structure Cu(tpp-*N*-O) with normal $S = 1/2$ spectra were obtained). The spectra are typical for planar copper(II) complex with $g_{\text{iso}} = 2.060$, $A_{\text{iso}}(^{63}\text{Cu}) = 79.7$ G, and $A_{\text{iso}}(^{14}\text{N}) = 13.0$ G for **3**-monomer in CHCl₃ (or CH₂Cl₂) at 20 °C (shown in Figure 4a) and with $g_{\parallel} = 2.19$ and $A_{\parallel}(^{63}\text{Cu}) = 182.1$ G for **3**-monomer in CHCl₃ (or CH₂Cl₂) at 77 K (shown in Figure 4b). These hyperfine couplings are similar in magnitudes to those of $g_{\text{iso}} = 2.082$, $A_{\text{iso}}(^{63}\text{Cu}) = 77$ G, and $A_{\text{iso}}(^{14}\text{N}) = 13.7$ G and of $g_{\parallel} = 2.196$ and $A_{\parallel}(^{63}\text{Cu}) = 166$ G obtained either from copper(II) octaethylporphyrin *N*-oxide [Cu(OEP-*N*-O)] **2** in toluene solution at 23 °C or from 1% of sample **2** doped in octaethylporphyrin *N*-oxide at the same temperature.² The pattern of nitrogen hyperfine coupling in the porphyrins *N*-oxide complex is significant, since in solution only seven lines with relative intensities of 1:3:6:7:6:3:1 are seen due to hyperfine interaction with nitrogen. Consequently, it appears that the unpaired spin in Cu(II) is interacting with three nearly equivalent nuclei rather than with four nitrogens, which is the usual case for porphyrins where a nine-line pattern is observed. We take this to be a further indication that this complex has a structure analogous to **3**-monomer shown in parts b and c of Figure 2. We cannot detect the resonances of $\Delta M_s = \pm 2$ at the

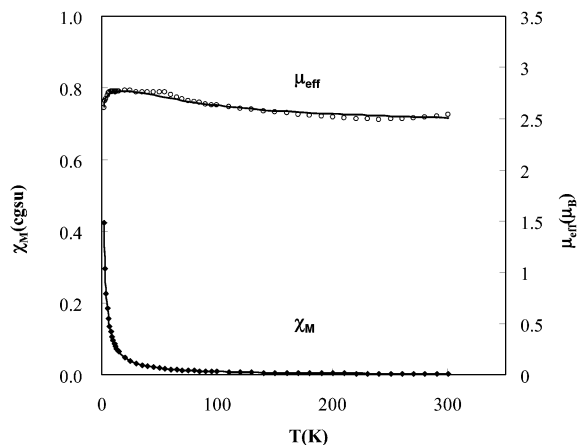


Figure 5. Temperature variation of the molar magnetic susceptibility (χ_m) and effective magnetic moment (μ_{eff}) for the powder sample of **3**-dimer in the range of 5–300 K. Points are the experimental data; solid lines represent the least-squares fit of the data to eq 8.

half-field in the parallel mode for **3**-dimer in CHCl₃ at 77 K; hence, there is no hint of forming a dimeric species in frozen solution at 77 K. This dimerization is the solid-state lattice effect of solid-state packing with an extremely small association constant.⁵²

Anisotropic Exchange Contribution to Zero-Field Splitting. Anisotropic exchange or pseudo-dipolar coupling (D_s^{ex}) has been widely used to explain the ground-state triplet zero-field splitting in copper dimer. Applying eq 7 and $r = 3.987$ Å for the Cu–Cu distance from X-ray data of **3**-dimer yields $D_s^{\text{dip}} = -0.058$ cm⁻¹ and this gives the opposite sign of the experimental ground-state zero-field splitting of $D_s = 0.074$ cm⁻¹. The values of D_s^{ex} ($= 0.132$ cm⁻¹) were obtained by substituting the values of D_s and D_s^{dip} into eq 6. These results show that in **3**-dimer the anisotropic exchange interaction D_s^{ex} gives the main contribution to D_s . Per the hypothesis that all of the D_s interactions arise due to dipolar interaction D_s^{dip} , the measured D_s values can be used to estimate the metal–metal (Cu²⁺...Cu²⁺) separation, r , by eq 7.⁵¹ Meanwhile, Bencini et al.³⁸ have emphasized that in the hypothesis of the dipolar interaction, the intensity of the $\Delta M_s = \pm 2$ transition in single crystals depends on r^{-6} . Moreover, the anisotropic exchange interaction D_s^{ex} gives the significant contribution to D_s in **3**-dimer. Hence, the approximate value of the Cu²⁺...Cu²⁺ distance from the half-field spectrum in **3**-dimer is not calculated.

Magnetic Properties. Magnetic data for complex **3**-dimer are reported in Figure 5 in the forms of χ_m and μ_{eff} versus T . As can be seen in Figure 5, the effective magnetic moment μ_{eff} for **3**-dimer gradually increases from 2.54 μ_B at 300 K to 2.77 μ_B at 10 K. This increase of μ_{eff} as the temperature decreases indicates that $S = 1$ is the ground state. Below 10 K, the μ_{eff} drops slowly to 2.60 μ_B at 2 K, reflecting the unavoidable zero-field splitting D_s of the ground state in **3**-dimer. The data were fitted by the Bleaney–Bowers^{22,53}

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(53) The influence of J interactions and zero-field effects on the energy levels of **3**-dimer is shown in Figure S2 in the Supporting Information.

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equation (8) and a term p which is fraction of any paramagnetic impurity, i.e., Cu^{2+} .

$$\chi_m = \frac{N g^2 \beta^2}{kT} \frac{2}{2 + e^{D_s/kT} + e^{(-2J + D_s/3)/kT}} (1 - p) + \frac{N p}{4} \frac{g^2 \beta^2}{kT} + \text{TIP} \quad (8)$$

Here, g is the average g value, TIP is the temperature-independent paramagnetism, N is Avogadro's number, p is the fraction of the paramagnetic impurity, and other symbols have their usual meanings. The best fits as represented in Figure 5 gave the values of $g = 2.01$, $2J = 68 \text{ cm}^{-1}$, $p = 0.06$, and the temperature-independent paramagnetism value $\text{TIP} = 1 \times 10^{-6} \text{ cm}^3/\text{mol}$ by substitution of the value of $D_s = 0.074 \text{ cm}^{-1}$ into eq 8. The negative value of $-2J$ indicates the ferromagnetic nature of spin coupling. This places the $S = 0$ state at about 68 cm^{-1} above the $S = 1$ ground state.

From eq 5, it is apparent that the hyperfine splitting $A_{\parallel,s}$ is half that for a related mononuclear copper ($S_1 = 1/2$) system (A_{\parallel}). This relation has been found to hold to a reasonably good approximation in several instances.^{39,40,54,55} The species **3**-monomer originating from the dissociation of **3**-dimer is a mononuclear Cu^{2+} ion, and its hyperfine constant is a qualified constant for A_{\parallel} . Although we found that $A_{\parallel,s}$ in **3**-dimer is $81 \pm 2 \text{ G}$ (Figure 3b), smaller than what would be expected from eq 5, since A_{\parallel} for the **3**-monomer is 182.1 G (Figure 4b). Moreover, the resulting copper hyperfine splitting in **3**-dimer expressed in wavenumbers with $A_{\parallel,s} =$

$(9.6 \pm 0.2) \times 10^{-3} \text{ cm}^{-1}$ is close to half of the corresponding hyperfine splitting in **3**-monomer with $1/2 A_{\parallel} = 9.35 \times 10^{-3} \text{ cm}^{-1}$. This supports that the magnetic exchange of copper pairs in **3**-dimer is in the strong coupling limit.

Conclusions

We have investigated two new metal complexes, namely, one paramagnetic, binuclear copper(II) and one diamagnetic, binuclear zinc(II) porphyrin *N*-oxide; complexes of **3**-dimer and **4**-dimer and their X-ray structures are established. We have also reported for the first time the ESR spectra of $[\text{Cu}(\text{tpp-}N\text{-O})_2]$ (**3**-dimer) in powder form at 293 and 4 K. Structural and ESR studies showed that **3**-dimer produces a spin triplet ground state formed by ferromagnetic coupling between the two $d^9 \text{ Cu}^{\text{II}}$ ions. The results clearly show that when pure single crystals of **3**-dimer are dissolved in CHCl_3 (or CH_2Cl_2) solvent, the principal copper species is **3**-monomer.

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Supporting Information Available: Figure S1 showing the atom numbering scheme for **4**-monomer; Figure S2 showing the influence of J interactions and zero-field effects on the energy levels of **3**-dimer; X-ray crystallographic files in CIF format for compounds **3**-dimer and **4**-dimer are also depicted. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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