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## 電子擴散系統的快速計算法（2／3）

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# Fast Algorithms for Electro-Diffusional Systems ${ }^{1}$ (Progress Report) 

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## 1 Background

Most biochemical processes involve macromolecules in solution. Usually, the macromolecule is surrounded by a hydration shell and is immersed in mobile ionic solvent. The corresponding electrostatic potential (in or around) is of central importance for understanding their functions [16]. Two typical approaches are the molecular dynamics simulation and the continuum modeling. The molecular dynamics simulation provides more adequate description at microscopic scale, but it requires a large number of adjustable parameters and computing resources. Indeed, it is impractical for the spatial and time scale we are interested in. On the other hand, the continuum approach is simple, adjustable parameter free, relatively inexpensive, and has been widely used [16].

The continuum model for molecules in solution was pioneered by Debye-Hückel in 1924 [2, 10]. The macromolecule such as protein or nucleotide represented by a structured and polarized clusters of charges. It sits inside a region $\Omega_{\text {in }}$ with low dielectric constant $\epsilon_{1}$, and is surrounded by a hydration shell $\Gamma$, which is immersed in a mobile ionic solvent in region $\Omega_{\text {out }}$ with high dielectric constant $\epsilon_{2}$. The hydration shell prevents the ionic solvent move into the inside region. Sometimes, the hydration shell could occupy a layer region (denoted by $\Omega_{\text {layer }}$ ) with dielectric constant $\epsilon_{2}$. In $\Omega_{\text {out }}$, there could be several kinds of ions. For easy presentation of our numerical model, we assume there are only two kinds of ions with opposite sign and same charge unit, and with total neutralized solution. The corresponding electrostatic potential $\phi$ satisfies the Poisson-Boltzmann equation $[2,18,19]$ :

$$
\begin{equation*}
-\nabla \cdot[\epsilon(x) \nabla \phi(x)]-\sum_{i} K(x) z_{i} \exp \left(-z_{i} \phi(x) / k T\right)=Q(x) . \tag{1.1}
\end{equation*}
$$

Here, $i$ is the index for the ith species of ions, $\epsilon(x)$ is the dielectric function taking value $\epsilon_{1}$ in $\Omega_{\text {in }}$ and $\epsilon_{2}$ in $\Omega_{\text {out }} \cup \Omega_{\text {layer }}, K(x)=0$ for $x \in \Omega_{\text {in }} \cup \Omega_{\text {layer }}$ and $K(x)=K$ for

[^0]$x \in \Omega_{\text {out }}$, a modified Debye-Hükel parameter, represents the ionic strength for solvent, the function
\[

$$
\begin{equation*}
Q(x)=4 \pi \sum_{k=1}^{m} q_{k} \delta\left(x-x_{k}\right) \tag{1.2}
\end{equation*}
$$

\]

represents the charge distribution of the macromolecule in $\Omega_{\mathrm{in}}$.
In the community of computational chemistry, there were decades of efforts of improvement on modeling by various groups of peoples such as Hornig (DelPhi package), McCammon, etc. Previous numerical models can be classified into grid-based methods such as finite difference, finite element and finite volume [7, 20, 13], and boundary integral methods such as boundary element for integrating the discretized linear systems such as SOR [7], multigrid [6] have also attracted attentions. The nonlinear counterpart was solved by direct iteration [20], or Newton's iteration [6]. However, most of above methods are low order, or not fast enough. Some of them indeed create severe numerical errors for large macromolecular such as DNA. For instance, one major error source comes from the treatment of singular point charge source. This point charge source term in a finite difference method is usually handled by direct discretization, or by distributing them into neighboring grid points. This creates $O\left(1 / h^{1}\right)$ error in $L^{\infty}$ near macromolecules. In addition, the treatment of discontinuous dielectric constant is also naive. The coefficient smoothing method commonly used in this community has $O(1)$ error in potential. The error may be larger for the force. These numerical problems leave room for mathematicians to do improvement.

## 2 What we have done so far

In these two years, we have done the following subprojects.

- "Accurate Evaluation of Electrostatics for Macromolecules in Solution," (with Jian-Guo Liu and Wei-Cheng Wang), Methods and Applications of Analysis, Vol. 10, No. 2, pp. 309 (2003). The main point is to use multipole method to handle the singular source and to use a body-fitting grid method to handle interface discontinuities.
- "New Formulation and Fast Poisson Solvers for Interface Problems in Polar coordinates," (with Zhilin Li, Wei-Cheng Wang and Ming-Chih Lai), The main thing addressed is to use a distance function to remove jump singularity of the source terms on interfaces. SIAM J. Sci. Comp., Vol. 25, No. 1, pp. 224-245 (2003).(SCI)
- "A spectral method for two-dimensional Poisson-Boltzmann equation," (with Chien-Chang Yen, Jann-Guo Liu and Ming-Chen Shiu), in preparation. The work here is to develop a spectral method when the interface is a circle or a ring.
- "A New Immersed Interface Method for Elliptic Equations with Discontinuous Ciefficients," (with Yu-Chen Shu), in preparation. In this research, we have develop a robust and simple second order finite method for elliptic equations with discontinuous coefficients. The underlying grid is regular, say Cartesian, and fixed. The interfaces may move and are represented by level sets. Interface problem is a well-konwn problem since the work of Peskin. The issue is to design a simple and high-order method under Cartesian grid and to allow interfaces dynamically move. We have compared our method with previous works of Peskin (immersed boundary method, see the review article [21]), LeVeque and Zhilin Li [25] (immersed interface method), Liu, Fedkiw and Kang [22] (ghost fluid method), our new method is the best. It is simple and second-order for both potentials and its gredients.


## 3 What we plan to do in the last year

### 3.1 The research subprojects

- We plan to write the paper for the new immersed interface method.
- We can do one of the following researches in the third year:
- We plan to extend the new immersed interface method to 3-d.
- We plan to implement the new immersed interface method to PoissonBoltzmann equations. The new element will be some nonlinear preconditioner.
- We plan to import some charge distribution data for some prototype protein from the protein data bank and to test our method for real drug design problems.


### 3.2 The difficulties and the approaches

- In 3-d, the difficulty needed to be resolved is the tangential derivative correction on the interfaces. Our 2-d method can be extended directly to 3 -d without difficulty. However, the algebraic multigrid method for solving the resulting linear system is needed to tune up.
- For solving Poisson-Boltzmann equation, we plan to find some nonlinear precondotioner to speed up the nonlinear iterations.
- The real problem is an implementation problem.


### 3.3 Expected results

After this work, we should have a method which is better than any current popular methods (Hornig, McCommon, Holst, etc. see references in the bibliography).

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