

行政院國家科學委員會專題研究計畫 成果報告

子計畫一：被動光環型網路之測試平台(1)

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行政院國家科學委員會國家型電信計畫成果報告

被動光環在都會網路之應用

子計畫一：被動光環型網路之測試平台()

計畫編號：93 - 2219 - E - 002 - 008 -

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1 中文摘要

多波長被動光環網路所需之測試節點需要許多獨立訊號源，以建立高速可程式網路互連測試平台，雖然商用高速設備可以購得但設備價格卻非常昂貴。本計畫專注於建立網路測試節點的重點在於設計並完成高速可程式網路互連測試平台，在這國家型電性計畫的各個階段過程中，此高速可程式網路互連測試平台可演進為一個乙太網路節點，並可處理每秒 10 億位元秒(10-Gb/s)的網路流量。目前此測試平台正使用硬體描述語言 (Verilog) 並在現場可程式邏輯陣列 (FPGA) 板上驗證，並配合高速多工/解多功晶片將資料作高低速轉換，任意資料模型產生以及錯誤計算的功能可以在現場可程式邏輯陣列上重新規劃執行資料處理。

除了測試平台外，具有高光纖色散及低成本特性的每秒 10 億位元秒(10-Gb/s)傳收機也已經完成。此傳收機建立在可直接調變的分布式回授雷射(DFB)，利用一個窄頻光濾波器可使傳輸距離達 60 公里，相較於傳統架構的 20 公里，確實得到很好的效果，應用在環狀半徑較小的被動式光環網路，低成本直調傳收機可用來取代高價的外部調變傳收機，使網路成本與架構更具彈性及競爭力。

Abstract:

The test-bed for multi-wavelength passive optical (POR) network requires many independent signal sources as high-speed programmable networking testing platform. Although commercial high-speed equipment is available, those equipments are typically very expensive. The focus of this project is to build a networking test-bed is the design and implement of the high-speed programmable networking testing platform. In the continuation of this project, the testing platform will evolve to a networking node capable to process 10-Gb/s Ethernet traffic. Currently, the testing platform is programmed in Verilog and verified in a FPGA test board. The FPGA is combined a Mux and DeMux circuits. The FPGA for pattern generator and error counting can be re-programmed for

data processing

In addition to the testing platform, low-cost high-dispersion-tolerance 10-Gb/s transceiver is also implemented. Based on directly-modulated of a simple distributed-feedback laser with a narrow-band optical filtering, the signal can be transmitted for a distance of 60-km with dispersion compensation as compared with a distance of only 20 km in conventional scheme. The low-cost transceiver can greatly lower the cost for POR with small diameter without the requirement of costly externally modulated transceiver.

2 研究成果內容

2.1 研究目的

A test-bed for multi-wavelength passive optical ring (POR) networks is constructed in this project for metro applications in dense cities. No active component is in the ring path to lower the cost of each node and the future operating and maintenance expenses. The POR networks are suitable for dense cities having communities (or residential complex) composed of many high-rise residential buildings or business centers with many office buildings, all of them in close proximity with other similar communities or business districts.

The testbed requires independent signal sources, low-cost high-dispersion-tolerance transceiver, and future advanced modulation scheme. Because of those requirement, we build the following:

1. Based on FPGA and Mux/DeMux ICs, some signal sources initially generate pseudo-random binary sequence (PRBS) to fill up the transmission links. Currently, the circuits is able to generate various types of PRBS.
2. Based on narrow-band filtered directly-modulated laser, low-cost optical transceivers are developed for a transmission distance of 60-km using low-cost directly modulated DFB laser.
3. Advanced phase-modulated optical signal is studied for metro and long-haul transmission systems. Differential phase-shift keying (DPSK) signal is studied under the influence of fiber nonlinearities, especially nonlinear phase noise.

2.2 測試平台設計

In building the test-bed, we require many independent signal sources to generate signal for testing propose. Fig. 1 shows a 10-Gb/s pattern generator that provide 10-Gb/s pseudo-random binary sequence (PRBS) to populate all WDM channels with

independent signals. In Fig. 1, the FPGA generates a 16-bit parallel PRBS with a predetermined timing offset between each other. When the parallel PRBS is combined using a 16:1 multiplexer, the multiplexed sequence is a true 10-Gb/s PRBS. We implement PRBS with sequence length of 2^7-1 , $2^{15}-1$, $2^{23}-1$, and $2^{31}-1$.

Also shown in Fig. 1, we also build the error detector to count the error of an input data with respect to the PRBS. The FPGA of the error detector shares the same FPGA of the pattern generator. With a 16:1 demultiplexer, the FPGA can implement a parallel error detector with fixed data rate the same as that of 10-Gb/s Ethernet WAN-PHY signals.

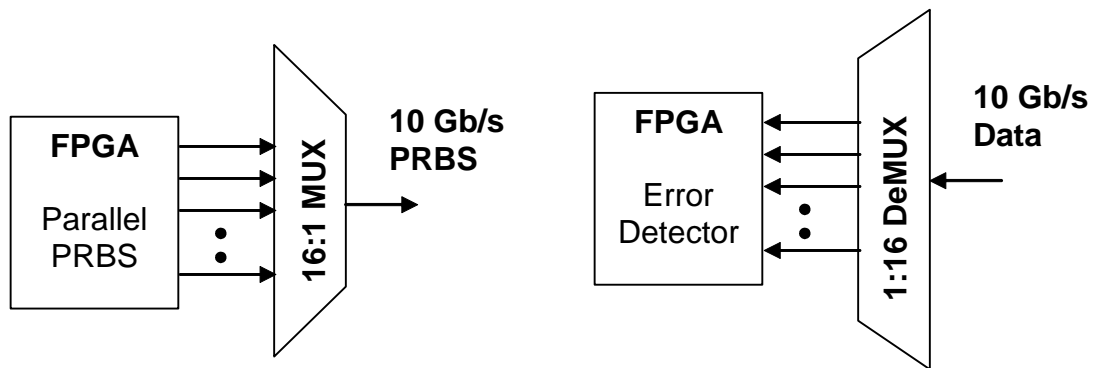


Fig. 1. 10-Gb/s Testing Platform.

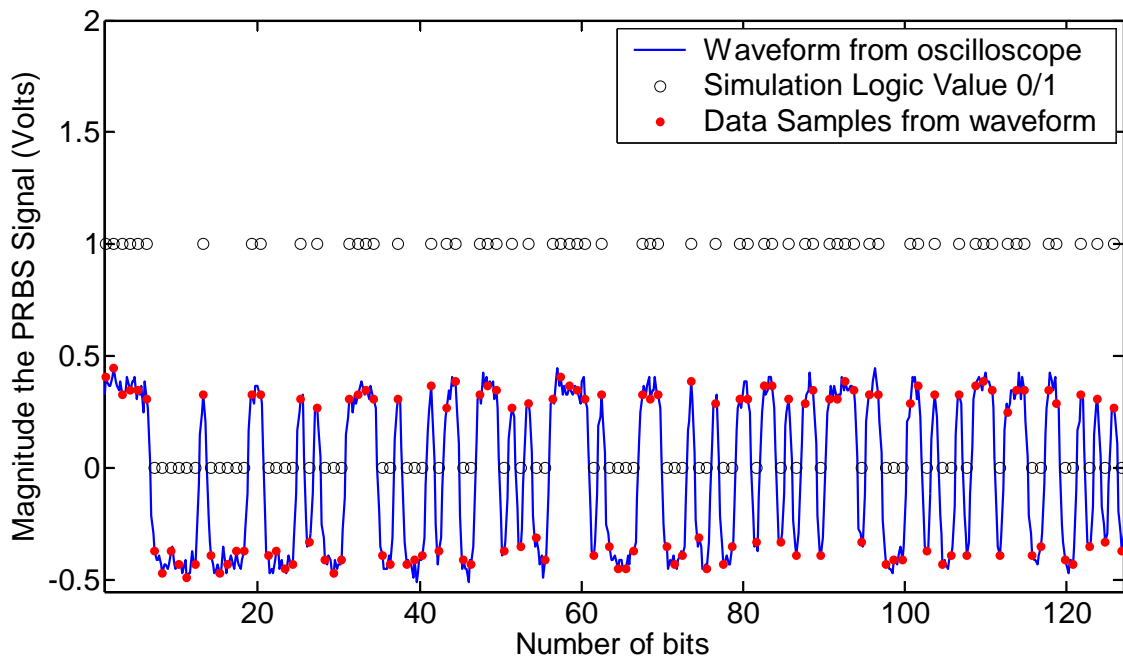


Fig. 2. FPGA-based PRBS generator signal verification for 2^7-1 pattern.

Currently, the logical functions of the PRBS are programmed in Verilog and verified in a FPGA test board. The test board is combined with a Mux and DeMux circuits. The FPGA for pattern generator and error counting can be re-programmed for data processing. In the future, the FPGA in each WDM node can emulate packet drop, pass, and add for the removal, by-passing, and insertion of Ethernet packet. We will implement a primitive node for Resilient Packet Ring according to IEEE 802.17 standard.

The FPGA based PRBS generator is built successfully. Fig. 2 shows the measured $2^7 - 1$ random pattern from an actual FPGA (blue solid line), data sampled from the wavelength (solid circle •) and the simulation using Verilog (circle o). The data rate

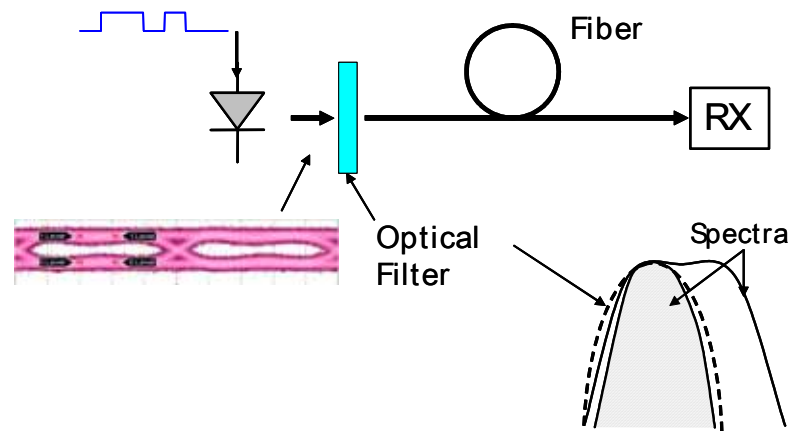


Fig. 3 Schematic for the experimental setup to generate an OOK signal based on FSK to OOK conversion. Also shown schematically the spectra before and after the optical filter. The eye-diagram shows the residual intensity modulation of the optical signal with an extinction ratio of 2 dB.

for Fig. 2 is actually 2 Gb/s, larger than the required speed of 622 Mb/s for Fig. 1. Although the pattern length can be larger than $2^7 - 1$, a short pattern is used as the oscilloscope can only capture a short pattern length.

2.3 低成本光收發機

A semiconductor diode laser by itself is a very simple and low-cost device. Light is generated when current is injected into the laser. In normal operation, the output power of the laser is monotonically increase with the injected current. On-off keying (OOK) signal can be generated by directly-modulated the diode laser with two different current levels. However, the emitted laser wavelength also depends on the injected current, resulting with dynamic frequency chirp that broadens the signal spectrum. When different portions of the signal are propagated with different speed

due to fiber chromatic dispersion, the transmission distance is limited, especially when low-loss 1.55- μm signals are used in the widely deployed standard single-mode fiber with a dispersion coefficient of $D = 17 \text{ ps/km/nm}$. Of course, the transmission distance can be increased when an external modulator is used to block or unblock the light. However, an expensive device, external modulator increases the cost of the system.

Conventionally, the usage of 1.55- μm directly modulated semiconductor laser limits the transmission distance to about 20 km, mostly due to the frequency chirp of the laser. Based on the same physical mechanism of frequency chirp, a diode laser can be directly frequency modulated to generate frequency-shift keying (FSK) signal. The FSK signal can be directly detected using an optical filter.

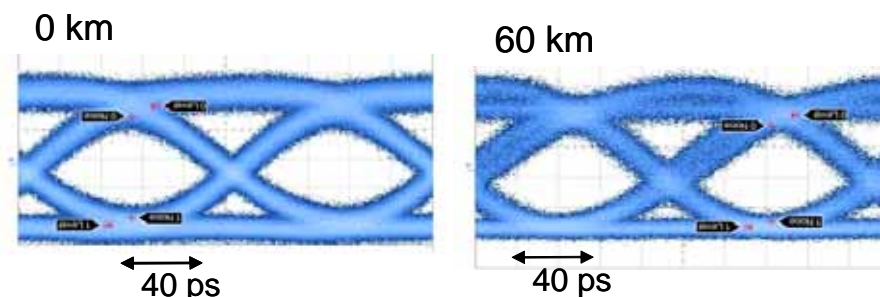


Fig. 4 Eye-diagram after 0 and 60 km of optical fiber. The eye-diagram for 0 km is for the converted OOK signal. After 60-km of optical fiber, the eye-diagram looks very similar.

When FSK signal is passing through an optical filter, the filter output signal is an OOK signal that can be directly detected using a conventional binary receiver. If FSK signal is transmitted through an optical fiber, the optical filter is located at the receiver. Fiber chromatic dispersion and optical filtering are both linear time-invariant systems for the optical signal. The performance of the system remains the same by exchanging the location of the fiber and the filter. Therefore, an OOK signal can be regenerated using an optical filter and then transmitted through the optical fiber.

Fig. 3 shows a semiconductor laser that generates a FSK signal through direct modulation. A passive optical filter converts the FSK signal to OOK signal that is transmitted through the optical fiber. In the design of the filter of Fig. 3, the notch of the filter should be located at the frequency corresponding to the lower intensity and the converted OOK signal has zero chirp at the “off” state.

The operation principle is further illustrated in Fig. 3 with the spectra before and after the passive optical filter. The operation principle is the same as most direct-

detection FSK systems but the optical filter that converts FSK to OOK signal is located at the transmitter instead of the receiver. Ideally for MSK, the optical filter should have a peak-to-notch bandwidth of 5-GHz, half the data rate, to give an OOK signal with infinite extinction ratio. In practice, the extinction ratio of the converted OOK signal is finite.

In Fig. 3, the laser is directly modulated with 10-Gb/s 2^7-1 pseudo-random binary sequence (PRBS) with the same operation parameters as Fig. 1(a). The optical filter has a 3-dB bandwidth of about 8.5 GHz, based on the monochromator in Agilent 86146B optical spectrum analyzer. Focused on the study of the dispersion tolerance, an optical amplifier is used between the semiconductor laser and the optical filter to compensate for the high optical filter loss (~ 7 dB). Following the optical filter is standard single-mode fiber with dispersion coefficient of about 17 ps/nm/km.

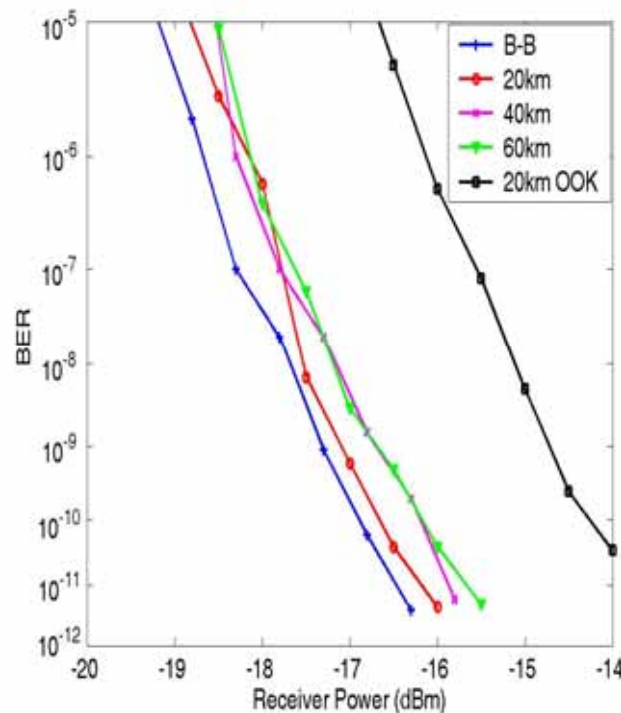


Fig. 5 BER as a function of received power for various length of optical fiber for 10-Gb/s signal.

In addition to direct frequency modulation, the FSK signal also has residual amplitude modulation and the corresponding eye-diagram is also shown in Fig. 3. The extinction ratio of the amplitude-modulated signal is about 2 dB. The extinction ratio of only 2-dB gives 6.5 dB of penalty for a standard OOK receiver.

Fig. 4 shows the eye-diagram of the converted OOK signal after passing through the

optical filter that has no difference with a regular OOK signal. The low-chirp OOK signal has an extinction ratio of about 8.5 dB. Fig. 4 also shows the eye-diagram after 60-km of optical fiber. The eye-diagram shows that there is little chromatic dispersion induced distortion to the signal. Even after 60-km of standard single-mode fiber, the eye-diagram does not show a big difference with that at 0-km of optical fiber.

Fig. 5 shows the measured bit-error-rate (BER) as a function of received power for the setup of Fig. 3 with 0, 20, 40, and 60 km of optical fiber. There is little chromatic dispersion induced penalty. Even for a distance up to 60 km, the maximum distance limited by our resource, the sensitivity penalty is less than 1 dB. Also shown in Fig. 3 is the BER after 20-km of optical fiber using the pure directly-modulated signal with a penalty of about 2 to 2.7 dB compared with the low-chirp signal. Even for just 20-km, the signal is up to 2 dB worse than the low-chirp signal with 60-km of optical fiber. Fig. 5 shows clearly the high dispersion tolerance of the OOK signal converted from a directly-modulated FSK signal.

We experimentally demonstrate a 10-Gb/s OOK transmitter based on directly modulated 1.55- μm semiconductor laser to generate FSK signal. A passive optical filter is used to convert the FSK signal to low-chirp OOK signal. The low-chirp OOK signal can be transmitted for a distance of 60-km of standard single-mode fiber with less than 1-dB of penalty induced by chromatic dispersion.

2.4 先進的調變型式

Recently, multilevel modulation signals like differential quadrature phase-shift keying (DQPSK) have received renewed attention to improve the spectral efficiency of a lightwave communication system. Not only for (D)QPSK signals, but also for the general class of quadrature-amplitude modulation (QAM) with and without differential operation, the transmitter of Fig. 6(a) is the usual method to generate a QAM signal using two Mach-Zehnder modulators (MZMs) within two $\pi/2$ phase difference paths of a Mach-Zehnder interferometer.

With two MZMs and an interferometer, the special modulator of Fig. 6(a) is difficult to fabricate. If the transmitter of Fig. 6(a) is implemented using discrete components of two MZMs within an interferometer, the transmitter is costly with many components. The transmitter in Fig. 6(a) also requires two bias controls for the MZMs and a phase control of the phase shifter. As shown subsequently, the QAM signal can also be generated using a single dual-drive MZM. When the technique is

applied to (D)QPSK signals, three different transmitters are invented with a drive signal having four, three, and two levels, respectively. The usage of only one dual-drive MZM greatly simplifies the design of the (D)QPSK transmitter.

Fig. 6(b) is the basic structure of a dual-drive MZM. The dualdrive MZM consists of two phase modulators that can be operated independently. The MZM can be fabricated using various materials, and LiNbO₃ is the most popular material. Almost all commercial long-haul dense-wavelength-division-multiplexed (DWDM) systems use LiNbO MZM. In the dual-drive structure of Fig. 6(b), the modulator chirp is adjustable

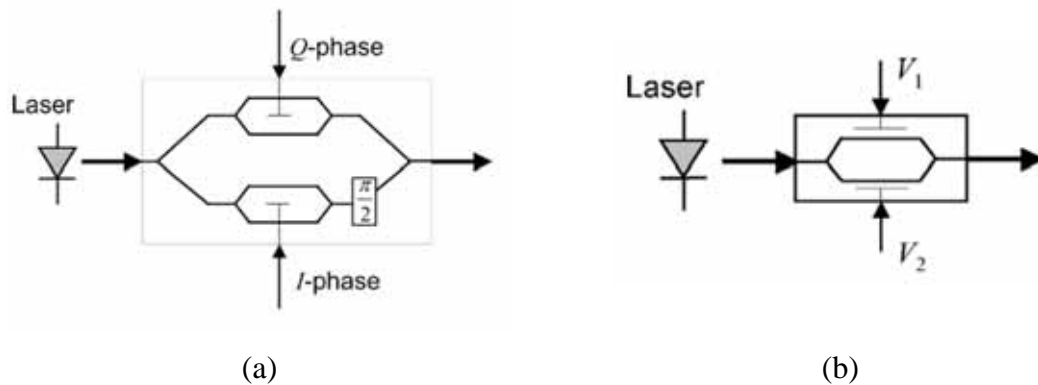


Fig. 6. Arbitrary quadrature signal transmitter based on (a) two single-drive MZMs in an interferometer and (b) one dual-drive MZM with two independent phase modulator in a Mach-Zehnder interferometer.

Fig. 7 show the eye diagram of the drive signal and the optical intensity between the RZ modulator and the (D)QPSK transmitter. Fig. 7(a) is the eye diagram when the conventional transmitter of Fig. 6(a) is used with two two-level drive signals having a peak-to-peak drive voltage of $2V_\pi$. Using the dual-drive transmitter of Fig. 6(b), the peak-to-peak drive voltage is reduced from $1.5V_\pi$ for a four-level signal to V_π for two- and three-level drive signals.

The output intensity of the conventional transmitter has optical intensity ripples between consecutive symbols. With two or three levels of drive signal, the output intensity of the dual-drive MZM also has ripples between consecutive symbols. The simplest two-level scheme of 7(c) has overshoot ripples doubling the output intensity. The ripple of the three-level signal of Fig. 7(d) with a dual-drive transmitter is similar to that of Fig. 7(a) with a conventional transmitter. If the RZ modulator is operated in the middle of the eye diagram of Fig. 7, the ripples are cancelled, and the

transmitter provides a constant pulse train. For an NRZ signal without an RZ modulator, overshoot ripples are equivalent to short optical pulses that are particularly detrimental and potentially give high signal distortion due to fiber nonlinearities. Without the RZ modulator, the two-level drive signals cannot be used for NRZ signals due to the overshoot ripples.

When the RZ modulator is used to convert the signal into a pulse train, the electrical drive signal of Fig. 7 must have a short rise–fall time such that the signal transition is outside the RZ pulses. Because of the overshoot ripples, the two-level transmitter of 7(c) requires the shortest rise–fall time from the electrical drive signal. If the rise–fall time is 20% the symbol interval, there is insignificant signal distortion due to the overshoot of Fig. 7(c). Even for a rise–fall time approaching 40% of the bit interval, the ripples in the RZ pulses are below half of the peak intensity.

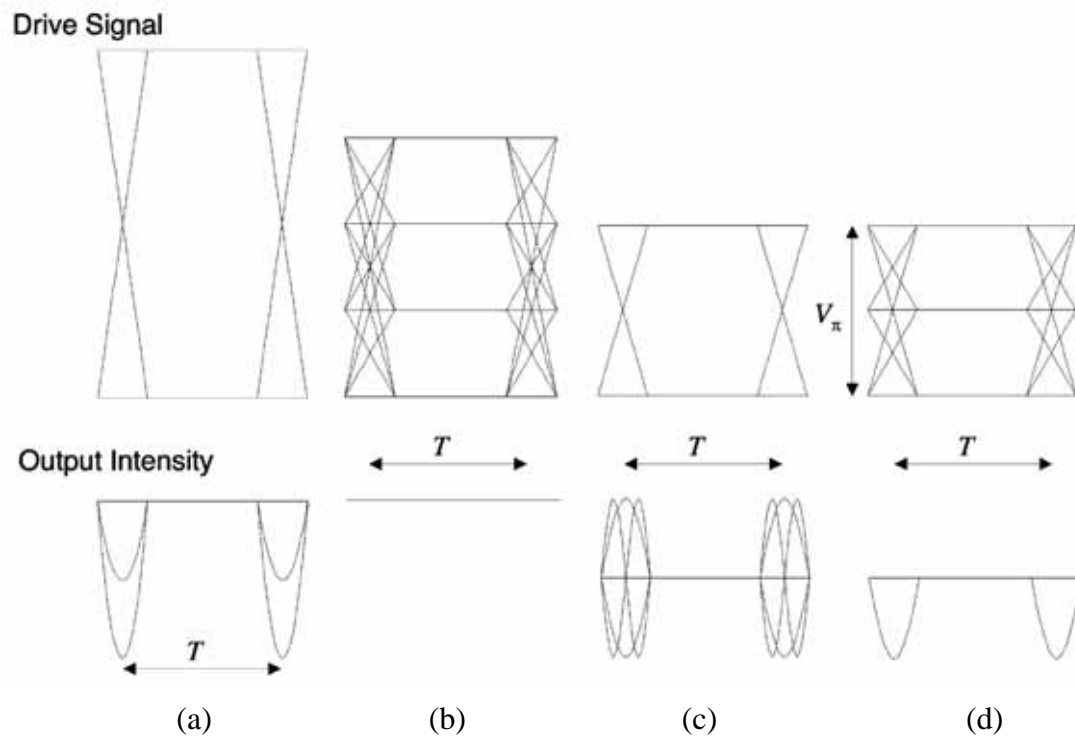


Fig. 7. Eye diagram of the drive signal and output intensity between two modulators of the RZ-DPSK transmitter when (a) the conventional transmitter of Fig. 1 and the dual-drive transmitter of Fig. 2 with (b) four-, (c) two-, and (d) three-level drive signals are used.

The ripples of the drive signal transfer to the optical signal. With the conventional transmitter of Fig. 6(a), no amplitude ripple of the drive signal transfers to the phase ripple. Even when the drive signal has a large ripple, the intensity ripple of the transmitted signal is compressed by the nonlinear transfer function of the MZM. For the (D)QPSK signal Fig. 7, the ripples from the drive signal may be increased by the

transmitter.

Regardless of the number of constellation points, all quadrature-amplitude modulation (QAM) signals can be generated using a single dual-drive Mach-Zehnder modulator. When the general method is applied to quadrature-phase-shift-keying (QPSK) signals, three different QPSK transmitters are shown with drive signals having four, three, or two levels. The usage of only one dual-drive modulator greatly simplifies the design of QAM and QPSK transmitters.

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4 計畫成果自評

本年度計畫執行符合本國家型電信計畫所預定完成之光環網路所需之硬體測試平台中每秒百億位(10-Gb/s)乙太廣域網路實體層測試訊號產生及接收器 [1]，利用 FPGA 及多工/解多工晶片及相關零組件設計一可程式化(Programmable)及可重置(Reconfigurable)硬體平台，具有可彈性設計的優點並可適用於多數每秒百億位(10-Gb/s)高速網路應用。

本研究可培養學生對系統整合以及高速電路設計訓練，此硬體平台尚可提供光通訊實驗測試誤碼率所需要的傳送與接收硬體設備，可訓練研究生不僅會用還能動手做出成品，藉此與理論驗證學以致用，具有高頻電路與系統整合實做經驗的研究生在未來職場上將炙手可熱，並提供業界所需的人才。

關於低成本的直調傳輸機架構，原本直調分布式回授雷射(DFB)的傳輸距離只能在 20 公里的範圍內，只能用在接取網路(Access Network)，利用價格低廉的窄頻光濾波器控制信號頻譜，即可將傳輸距離提升至 60 公里而可運用在小範圍的都會網路 [2],[3]，對於建置成本與網路彈性功能接鑄銖必較的產業界，將可提升價格及技術的競爭力。

本計畫除務實解決業界需求，在相關理論研究也不落人後，對於高速差分相位調變(DPSK)訊號收光纖非線性效應的研究，我們首先研究單一通道內訊號間的四波混合現象(Intra-channel Four-Wave Mixing, IFWM)，並比較其與非線性相位雜訊對 DPSK 訊號的影響，此研究榮膺 STFOC`05 最佳論文獎(詳見附錄) [4]，研究速率可調系統則獲 STFOC`05 傑出學生論文獎(詳見附錄) [5]，此外還有兩篇研究在 STFOC`05 發表 [3],[6]。

我們秉持著審慎的態度繼續研究 DPSK 信號在傳輸系統中受非線性效應影響的分析，用嚴謹的方法分析 RZ-DPSK 信號在分波多工系統中因為通道上信號非週期性特性所造成的相位交互影響調變效應(XPM)，此一研究(詳見附錄)在 2005 年 7 月發表在亞洲重要的光通訊研討會-OptoElectronics and Communication Conference (OECC) –舉辦地點在韓國首爾，由博士班研究生上台發表論文 [7]，此外尚有 3 位碩士班學生所研究的題目發表篇壁報論文 [2],[8],[9]。

本計畫主持人何鏡波副教授也受 OECC`05 大會邀請演講，發表 DPSK 信號在光傳輸系統非線性效應的統計特性的研究成果 [10]，此外在計畫經費補助下，尚參與香港及德國舉辦的國際會議 [11], [12]，對台灣在光通訊的能見度將有所幫助。在國際期刊方面，關於單一通道四波混合對 DPSK 信號的研究被 IEEE Photonics Technology Letters 所接受 [13],[14]，此外我們提出新的方式只需一個雙埠驅動的外部調變器(Dual-drive MZM)即可產生 DQPSK 信號，可降低訊號產生的成本，此論文發表於 IEEE/OSA Journal of Lightwave Technology [15]，對光通訊通道容量研究以及非線性效應的補償也有很好的成果 [16],[17]，總計本計畫經費支援下產生一篇碩士論文 11 篇國際會議論文以及 5 篇國際期刊論文。

XPM-Induced Crosstalk for RZ-DPSK Signals in Highly Dispersive Transmission Systems

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Abstract

In highly dispersive RZ-DPSK WDM systems, the overlap of optical pulses gives non-periodic optical intensity variation that gives phase noise through cross-phase modulation (XPM). The variance of the phase noise is evaluated analytically.

1 Introduction

Return-to-zero differential phase-shift keying (RZ-DPSK) signal has widely studied recently for long-haul transmission [1]. With low peak power, DPSK signal can tolerate larger fiber nonlinearities than on-off keying. If the RZ-DPSK signal maintains as a periodic pulse train in its intensity profile, cross-phase modulation (XPM) induces the same phase shift to all optical pulses in a wavelength-division-multiplexed (WDM) system [2]. However, with pulse broadening in dispersive transmission, overlapped pulses induce non-periodic intensity profile. This paper provides a method to calculate XPM-induced phase noise from pulse overlap.

2 XPM from Overlap RZ-DPSK Pulses

The XPM-induced periodic phase shift from periodic intensity does not degrade DPSK signals. RZ-DPSK signal is typically launched as non-overlap periodic pulse train at the transmitter. However, after a short distance, the pulse is broadened by fiber dispersion that gives pulse overlap. The RZ-DPSK signal has non-periodic intensity profile. The non-periodic intensity gives XPM-induced phase variation that degrades the performance of DPSK signals.

If the DPSK signal has a pulse shape of $g(z, t)$ in electric field as a function of distance z and time t , the two adjacent pulses give intensity profile of $|g(z, t) \pm g(z, t-T)|^2$. If the adjacent pulses of $g(z, t)$ and $g(z, t-T)$ are overlap with each other, there is an equivalent

non-periodic modulation of $\pm 2\Re\{g(z, t)g^*(z, t-T)\}$, where $\Re\{\cdot\}$ is the real part of a complex number. Of course, the equivalent modulation is important only if the pulses of $g(z, t)$ and $g(z, t-T)$ overlap with each other. Combined with XPM, this equivalent modulation gives non-periodic phase to adjacent WDM channels.

When previous papers [3-4] studied the impact of XPM to DPSK signal, only signal and noise beating is included and pulse-overlap is not investigated. The model here follows the method of [3] in which periodic intensity does not degrade DPSK signals.

Following most analysis of XPM in [2-3], we assume that a Gaussian pulse of $g(t) = A_0 \exp[-(t/\tau_0)^2/2]$ in electric field is launched into the fiber to transmit RZ-DPSK signal, where $A_0 > 0$ is the peak amplitude and τ_0 is the initial $1/e$ -pulse width. The optical power of the RZ-DPSK signal is $P_0 = \pi A_0^2 \tau_0 / T$.

After a fiber distance of z , the optical Gaussian pulse in term of electric field becomes

$$g(z, t) = \frac{A_0}{\sqrt{1 - j\beta_2 z / \tau_0^2}} \exp\left[-\frac{1}{2} \times \frac{t^2}{\tau_0^2 - j\beta_2 z}\right], \quad (1)$$

where β_2 is the group-velocity dispersion coefficient of fiber. The $1/e$ -pulse width of $g(z, t)$ is $(\tau_0^2 + A_0^2 z^2 / \tau_0^2)^{1/2}$.

Assuming RZ-DPSK signal of $A(z, t) = \sum_k b_k g(z, t - kT)$ at distance of z , where $b_k = \pm 1$ is the phase modulated to the optical pulses. The overall amount of XPM-induced phase shift is equal to

$$\phi(t) = 2\gamma \int_0^L |A(z, t - d_\omega z)|^2 \otimes h_z(t) e^{-\alpha z} dz, \quad (2)$$

where \otimes denotes convolution and $h_z(t)$ is the impulse response of $H_z(\omega) = \cos(\beta_2 z \omega^2 / 2)$ due to fiber dispersion [5], d_ω is the walk-off parameter, and α is the fiber attenuation coefficient. Similar to a digital modulated

signal, the phase of (2) can be expressed as

$$\phi(t) = 2\gamma \sum_{k=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} b_k b_{k+m} q_m(t - kT) \quad (3)$$

where

$$q_m(t) = \int_0^L [g(z, t - d_\omega z) g^*(z, t - d_\omega z - mT)] \otimes h_z(t) e^{-\alpha z} dz \quad (4)$$

The power spectral density of $\phi(t)$ is evaluated analytically similar to digital modulated signal. The differential phase that degrades a DPSK signal is equal to $\Delta\phi = \phi(t) - \phi(t-T)$ has a variance of

$$\sigma_{\Delta\phi}^2 = 4 \int_{-1/T}^{1/T} \Phi_\phi(2\pi f) \sin^2(\pi f T) df, \quad (5)$$

where $\Phi_\phi(2\pi f)$ is the power spectral density of $\phi(t)$ and the integration is between the data bandwidth of $\pm 1/T$, also the first lobe of $\sin^2(\pi f T)$.

3 Numerical Results

Fig. 1 shows the normalized phase shift standard deviation (STD) of $\sigma_{\Delta\phi}/\langle\Phi_{NL}\rangle$ as a function of channel separation of $\Delta\lambda$ for both medium and high dispersion coefficient of 3.5 and 17 ps/km/nm, where $\langle\Phi_{NL}\rangle$ is the mean nonlinear phase shift. The initial pulse width of the Gaussian pulses is equal either $\tau_0 = 2.5$ and 5 ps.

The STD of $\sigma_{\Delta\phi}$ decreases with the increase of channel separation. The STD of $\sigma_{\Delta\phi}$ for short pulse is significantly smaller than long pulse when channel separation is larger than 1.6 nm. The fiber link of Fig. 1 has a loss coefficient of $\alpha = 0.2$ dB/km for a distance of $L = 100$ km. The mean nonlinear phase shift is $\langle\Phi_{NL}\rangle = \gamma P_0 L_{\text{eff}}$, where $L_{\text{eff}} = [1 - \exp(-\alpha L)]/\alpha$ is the effective nonlinear length, γ is the fiber nonlinear coefficient, and P_0 is the launched power.

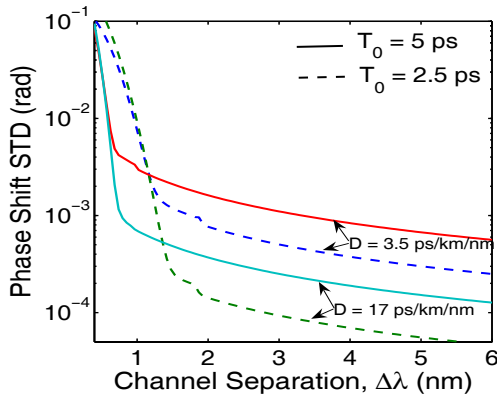


Fig. 1 The normalized phase shift STD of $\sigma_{\Delta\phi}/\langle\Phi_{NL}\rangle$ versus channel separation of $\Delta\lambda$ with $D = 3.5$ and 17 ps/km/nm.

Fig. 2 shows the normalized phase shift STD of $\sigma_{\Delta\phi}/\langle\Phi_{NL}\rangle$ as a function of fiber dispersion coefficient of D . The channel separation is $\Delta\lambda = 0.8$ nm for $\tau_0 = 5$ ps and $\Delta\lambda = 1.6$ nm for $\tau_0 = 2.5$ ps, corresponding to a frequency separation of 100 and 200 GHz at the wavelength around 1.55 μm .

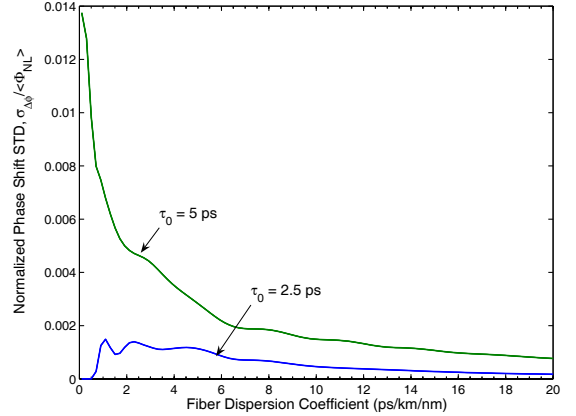


Fig. 2 The normalized phase shift STD of $\sigma_{\Delta\phi}/\langle\Phi_{NL}\rangle$ as a function of fiber dispersion coefficient of D .

Fig. 2 shows that pulse overlap induced phase noise depends strongly on the initial pulse width of τ_0 . For long initial pulse width, the initial pulse overlap dominates the pulse overlap and the STD of $\sigma_{\Delta\phi}$ decrease with increase of fiber dispersion. For short initial pulse, there is no initial pulse overlap. For short optical pulse without initial pulse overlap, the phase noise STD initially increases with fiber dispersion and gives a peak at dispersion coefficient 2 to 6 ps/km/nm and then decreases with the further increases of fiber dispersion.

4 Conclusion

The XPM-induced phase noise due to overlapped pulses is studied analytically the first time. The phase noise STD decreases with channel separation and increases with initial pulse width. The phase noise STD is also commonly decreases with the increase of fiber chromatic dispersion.

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Comparison of Nonlinear Phase Noise and Intrachannel Four-Wave-Mixing in RZ-DPSK Systems

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Abstract — Self-phase modulation (SPM)-induced nonlinear phase noise is reduced with the increase of fiber dispersion but intrachannel four-wave-mixing (IFWM) induced ghost pulses are increased with dispersion. The standard deviation of nonlinear phase noise is about three times that from IFWM.

1. INTRODUCTION

Self-phase modulation (SPM) induces nonlinear phase noise [1], adds directly to the signal phase, is the major degradation for RZ-DPSK signals. When RZ pulse broadens by chromatic dispersion and overlaps with each other, the pulse-to-pulse interaction gives intrachannel cross-phase modulation (IXPM) and four-wave-mixing (IFWM). While IXPM has no effect on DPSK signals, IFWM adds ghost pulses to each DPSK RZ pulse [2-4].

2. NONLINEAR PHASE NOISE VARIANCE

For a comparison to IFWM, nonlinear phase noise is evaluated based on the model of [2-4]. Assumed a Gaussian pulse with an initial $1/e$ pulse width of T_0 , the k th pulse along the fiber is

$$u_k(z, t) = \frac{A_k T_0}{(T_0^2 - j\beta_2 z)^{1/2}} \exp\left[-\frac{(t - kT)^2}{2(T_0^2 - j\beta_2 z)}\right] \quad (1)$$

where $A_m = \pm A_0$ is the pulse amplitude modulated by either 0 or π phases, β_2 is the coefficient of group velocity dispersion, and T is the bit interval. The constant factor of fiber loss is ignored in (1) but includes afterward. The overall ghost pulse is equal to

$$j\gamma \int_0^L \left[u_k(z, t) u_l(z, t) u_m^*(z, t) \right] \otimes h_{-z}(t) e^{-\alpha z} dz, \quad (2)$$

where \otimes denotes convolution, and L is the fiber length, and α is the fiber attenuation coefficient. The impulse response of $h_{-z}(t)$ provides dispersion compensation for $h_z(t)$, the fiber chromatic dispersion.

To be consistent with the model for IFWM of (2), for the pulse of $u_0(z, t)$, the SPM-induced nonlinear force including amplifier noise of $n(z, t)$ is equal to

$$j\gamma \left[u_0(z, t) + n(z, t) \right] \left| u_0(z, t) + n(z, t) \right|^2. \quad (3)$$

For the signal, nonlinear force is $j\gamma u_0 |u_0|^2$ or that of (2) with $k = l = m = 0$. The nonlinear force associated with nonlinear phase noise has two different terms of $2j\gamma |u_0(z, t)|^2 n(z, t)$ and $2j\gamma u_0^2(z, t) n^*(z, t)$ when all quadratic or higher-order terms of the noise are ignored.

With white input noise $n(0, t)$ has a variance of $E\{n(0, t+\tau)n^*(0, t)\} = 2\sigma_n^2 \delta(\tau)$, the fiber dispersion is included as $n(z, t) = n(0, t) \otimes h_z(t)$ and $E\{n(z, t+\tau)n^*(z, t)\} = 2\sigma_n^2 \delta(\tau)$, where σ_n^2 is the noise density per dimension. For $2j\gamma |u_0(z, t)|^2 n(z, t)$, the nonlinear noise corresponding to (2) is equal to

$$\Delta u_n(t) = 2j\gamma \int_0^L \left[|u_0(z, t)|^2 n(z, t) \right] \otimes h_{-z}(t) e^{-\alpha z} dz. \quad (4)$$

The temporal profile of $\Delta u_n(t)$ can be represented by the variance of $\Delta u_n(t)$, $\sigma_{\Delta u_n(t)}^2 = E\{|\Delta u_n(t)|^2\}$, as a function of time. Similar to $2j\gamma |u_0(z, t)|^2 n(z, t)$, the $2j\gamma u_0^2(z, t) n^*(z, t)$ has a temporal profile of $\Delta u_n(t)$ contributed to the nonlinear phase noise.

Fig.1 shows the standard deviation (STD) of $\Delta u_n(t)$ and $\Delta u_n(t)$ for typical fiber dispersion coefficients of $D = 17$ and 3.5 ps/km/nm. The initial launched pulse has an $1/e$ width of $T_0 = 5$ ps. The fiber link is $L = 100$ km with attenuation coefficient of $\alpha = 0.2$ dB/km. Fig.1 shows that the nonlinear force of $\Delta u_n(t)$ due to the beating of $|u_0(z, t)|^2$ with $n(z, t)$ is far larger than that of $\Delta u_n(t)$ due to the

附錄二

beating of $u_0^2(z,t)$ with $n^*(z,t)$. In term of power, the variance of $\Delta u_n(t)$ is about 1% of the variance of $\Delta u_n(t)$.

If the nonlinear force of $\Delta u_n(t)$ is passing through an optical filter with an impulse response of $h_o(t)$, the filter output at the time of mT is

$$\zeta_{0,m} = \int_{-\infty}^{+\infty} h_o(mT-t) \Delta u_n(t) dt. \quad (5)$$

The SPM phase noise from $\zeta_{0,0}$ is the noise generated by the beating of $|u_0(z,t)|^2$ with $n(z,t)$ and affect the DPSK pulse at $t = 0$. The term of $\zeta_{1,1}$ is IXPM phase noise from the beating of $|u_0(z,t)|^2$ with $n(z,t)$ and affect the DPSK pulse at $t = T$. Due to IXPM, the DPSK pulse at $t = 0$ also affects by the beating of $|u_1(z,t)|^2$ (the pulse at $t = T$) with $n(z,t)$ to give the IXPM phase noise of $\zeta_{1,0}$.

Followed the model of [4], the differential nonlinear phase noise from both SPM and IXPM phase noise is,

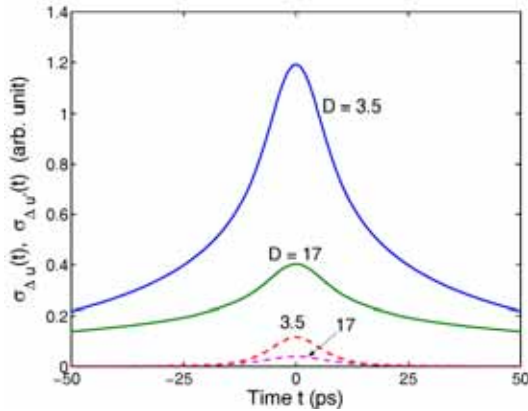


Fig. 1. The temporal distribution of nonlinear force due to the beating of signal with noise. The solid lines and dash line are for $\Delta u_n(t)$ and $\Delta u_n(t)$, respectively.

$$\delta\phi_n = \frac{1}{A_0} \Im \left\{ \sum_m \zeta_{m,0} \right\} - \frac{1}{A_1} \Im \left\{ \sum_m \zeta_{m,1} \right\} \quad (6)$$

where $\Im\{\cdot\}$ denotes the imaginary part of a complex number. For simplicity, $A_0 = A_1$ is assumed for the same transmitted phase in consecutive symbols. Using the property that the real and imaginary parts of $\zeta_{k,m}$ are independent and identically distributed, the variance of $\delta\phi_n$ is

$$\sigma_{\delta\phi_n}^2 = \frac{1}{A_0^2} \sum_{m_1} \sum_{m_2} \left(E \left\{ \zeta_{m_1,0} \zeta_{m_2,0}^* \right\} - E \left\{ \zeta_{m_1,0} \zeta_{m_2,1}^* \right\} \right). \quad (7)$$

For an N -span system, the amplifier noise at the first span is the smallest and that in the last span is the largest. From [1],[5], for large number of fiber spans with the identical span repeated one after another, the overall phase noise variance is $\sigma_{\Delta\phi_n}^2 \approx N^3 \sigma_{\delta\phi_n}^2 / 3$. The signal-to-noise ratio (SNR) is defined as $\rho_s = \pi^{1/2} T_0 |A_0|^2 / (2 N \sigma_n^2)$. The

mean nonlinear phase shift is $\langle \Phi_{NL} \rangle = N \gamma L_{\text{eff}} P_0$ where P_0 is the launched power and $L_{\text{eff}} = (1 - e^{-\alpha L}) / \alpha$ is the effective fiber length. The variance of nonlinear phase noise of $\sigma_{\Delta\phi_n}^2$ is proportional to $\langle \Phi_{NL} \rangle / \rho_s$, similar to that in [1],[5].

Fig. 2 shows the phase noise STD of $\sigma_{\Delta\phi_n}^2$ as a function of the fiber dispersion coefficient of the fiber link with Gaussian optical filter has a transfer function of $H_o(\omega) = \sqrt{1 + t_o^2 / T_0^2} \exp(-t_o^2 \omega^2 / 2)$. The system has a mean nonlinear phase shift of $\langle \Phi_{NL} \rangle = 1$ rad and SNR of $\rho_s = 20$ (13dB), corresponding to an error probability of 10^{-9} for DPSK signal with only amplifier noise.

The same as Fig. 1 with $T_0 = 5$ ps, Fig. 2 further assumes 40-Gb/s systems with $T = 25$ ps and an optical matched filter of $t_o = 5$ ps because larger bandwidth increases the nonlinear phase noise. Fig. 2. also shows the corresponding phase STD due to IFWM calculated by the method of [4].

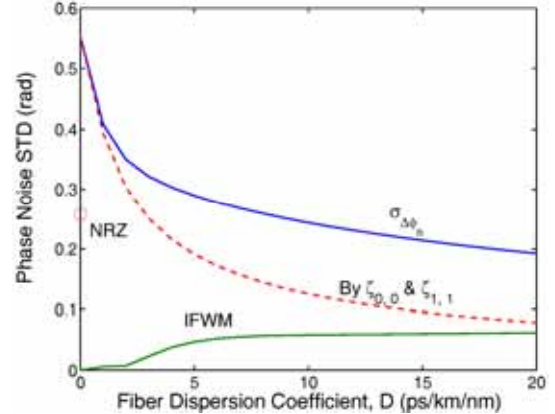


Fig. 2. The phase noise STD due to nonlinear phase noise and IFWM. The dashed-line is SPM phase noise from $\zeta_{0,0}$ and $\zeta_{1,1}$ alone.

For Gaussian optical filter, the variance of $\sigma_{\Delta\phi_n}^2$ can be derived analytically. Other types of optical filter may need another layer of integration.

4. CONCLUSION

For an initial pulse width of $T_0 = 5$ ps, the phase noise STD from nonlinear phase noise is about three times larger than that from IFWM at large fiber dispersion of $D = 17$ ps/km/nm. Nonlinear phase noise typically degrades a DPSK signal more than IFWM ghost pulses.

附錄二

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On-Off Keying with Different Mark-Ratio for Multi-Rate Transmission

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Abstract — A dynamic reconfigurable multi-rate transceiver is designed based on unequally probable on-off keying, or equivalently with different mark-ratio. Experimental measurement shows that the receiver sensitivity is improved by 4-6 dB by the halving the mark ratio.

1. INTRODUCTION

In a dynamic reconfigurable optical network, a transceiver may need to adapt to different data rates due to different requirement from customer or limitation in the link budget. If a transceiver is originally operated in 10-Gb/s, after reconfiguration, the new fiber link may have excessive fiber loss and the transceiver must lower its speed to 2.5 Gb/s requiring less receiver power or less effective signal-to-noise ratio.

Focused on the design of electronic circuit to detect the data rate and recover the data, conventionally, a multi-rate transceiver is mostly operated in low-speed from 52 Mb/s to 1.3 Gb/s [1-2]. The receiver sensitivity improvement from high to low data rates is not the major concerns in those works. Here, a new multi-rate transceiver is demonstrated based on unequally probable on-off keying, similar to change mark ratio of the data stream.

The multi-rate transceiver is operated with a maximum data rate of 10 Gb/s. In the multi-rate transceiver, the receiver sensitivity is inversely proportional to the data rate. If the data rate is reduced to 5 Gb/s, for example, half of the bits in predefined locations are always at the off state, saving 50% of the power. Ideally, for thermal-noise dominant receiver, the receiver sensitivity improves by 6 dB, providing further system margin or link loss budget. The analog front-end of both transmitter and receiver remains the same but the digital post-processing is changed for multi-rate operation.

2. PRINCIPLE OF UNEQUALLY PROBABLE ON-OFF KEYING

In a simplified analysis, the BER of an on-off keying receiver is given by the Q factor of

$$Q = (I_1 - I_0) / (\sigma_1 + \sigma_0), \quad (1)$$

where I_1 and σ_1 are the average photocurrent and noise variance for the on state, and I_0 and σ_0 are that for the off-state. With unequal probability of p_1 and p_0 , the average photocurrent is $I_s = I_1 \times p_1 + I_0 \times p_0$ and $p_1 + p_0 = 1$. With infinite extinction ratio of $I_0 = 0$, we obtain $I_1 = I_s / p_1$ and the Q factor is equal to $I_s / p_1 / (\sigma_1 + \sigma_0)$.

First of all, unequally probable on-off keying changes the effective data rate of the signal. For a nominally 10-Gb/s transceiver, if the signal has a mark ratio of $1/4$ instead of $1/2$, the effective data rate is just 5 Gb/s. In the parallel to serial multiplexing of digital data, some lanes can remained zero all the time. In the serial to parallel demultiplexing, the signal processing afterward can always skip certain lanes with all-zero data stream. The analog front-end of clock-and-data recovery (CDR), and serializer and deserializer (SerDes) remain the same. The CDR can operate for the full-rate 10-Gb/s data stream, or other effective data rates with different mark-ratio.

For system dominated by thermal noise or spontaneous-spontaneous beat noise, the noise variance of σ_1 and σ_0 are both independent of the photocurrents of both I_0 and I_1 . The change of p_1 from $1/2$ to $1/4$ doubles the Q factor and provides 6-dB sensitivity improvement. For system dominated by signal-spontaneous noise, the on-state noise of σ_1^2 is proportional to I_1 . Ignore the effect of σ_0 , the Q factor is proportional to $(I_s / p_1)^{1/2}$. The change of p_1 from $1/2$ to $1/4$ provides 3-dB sensitivity improvement. Depending on system configuration, the improvement of changing p_1 from $1/2$ to $1/4$ provides 3- to 6-dB sensitivity improvement.

5. EXPERIMENTAL VERIFICATION

Fig. 1 shows the experimental setup for the measurement of the BER of the multi-rate signaling at different transmission rates. A pulse pattern generator

附錄三

(PPG) is used to generate a 10 Gb/s $2^{31}-1$ PRBS signal. The PRBS is amplified by a driver amplifier to drive a Mach-Zehnder modulator (MZM) used with a laser source having the wavelength of 1553.28 nm. An optical attenuator is used to simulate fiber loss. The receiver is assumed as a preamplifier receiver with optical gain provided by a low-noise EDFA. Followed the EDFA is a second optical attenuator to maintain the same optical power to the receiver at different data rate, mainly for comparison purpose. Ideally, the second attenuator is not required but is used here to provide further flexibility. The receiver is followed by the error detector (ED).

At first, PPG generates a PRBS with mark ratio of $\frac{1}{2}$ for an effective data rate of 10 Gb/s. Fig. 2a shows the corresponding eye-diagram. Fig. 3 measured the bit-error-rate (BER) as a function of optical receiver power of the optical amplifier. Afterward, the PPG generates a PRBS with a mark ratio of $\frac{1}{4}$ for an effective data rate of 5 Gb/s. Fig. 2b also shows the corresponding eye-diagram. The BER as a function of received power is also measured and shown in Fig. 2. Finally, the PPG generates the data with a mark ratio of $\frac{1}{8}$ for an effective data rate of 2.5 Gb/s and an eye-diagram of Fig. 2c.

The required received power for a BER of 10^{-9} is -26.9 dBm for 10-Gb/s signal. The required received power for the effectively 5-Gb/s signal is -32.7 dBm and that for 2.5 Gb/s signal is -36.7 dBm. Comparing the BER as a function of received power in Fig. 3, 10-Gb/s operation of the multi-rate transceiver requires about 5.8 dB larger received power than 5-Gb/s and 5-Gb/s operation requires 4.0-dB larger received power than 2.5 Gb/s operation. Fig. 3 clearly shows the advantage of the multi-rate transceiver to provide adaptive data-rate with different receiver sensitivity.

In practice, the multi-rate transceiver can be operated with better granularity than that in Figs. 2 and 3. If the operation depending on the number of parallel data in the SerDes, the multi-rate transceiver can be operated for an increment of $\frac{1}{16}$ the data rate in the standard SFI-4 interface [3]. Of course, the implementation of the multi-rate transceiver does not need to conform to a standard interface and better granularity can be provided.

6. CONCLUSION

Unequally probable on-off keying is used for a dynamic reconfigurable multi-rate transceiver. Experiment shows that 10-Gb/s operation of the multi-rate transceiver requires about 5.8 dB larger received power than 5-Gb/s and 5-Gb/s operation requires 4.0 dB larger received power than 2.5 Gb/s operation.

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附錄三

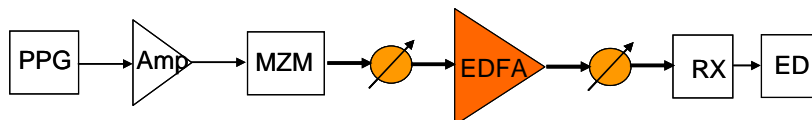


Fig. 1 The schematic diagram of experimental setup.

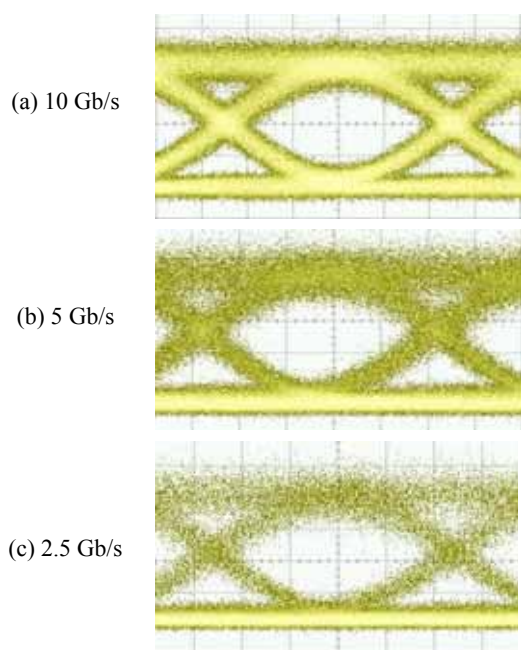


Fig. 2 Measured eye-diagram.

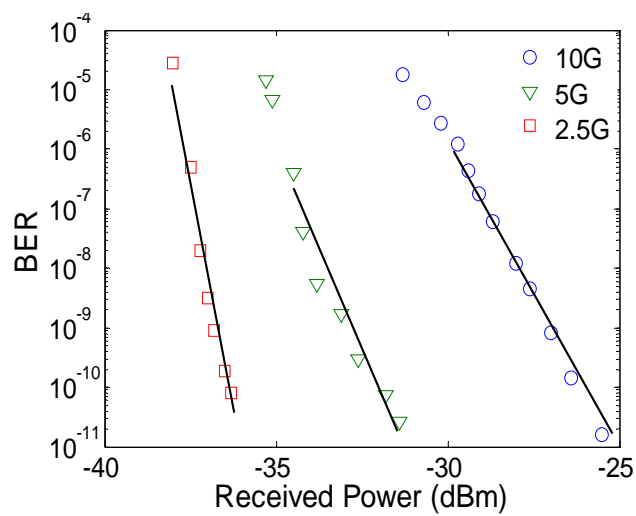


Fig. 3 Measured BER as a function of received power.